

回転大質量星の重力崩壊と爆発現象

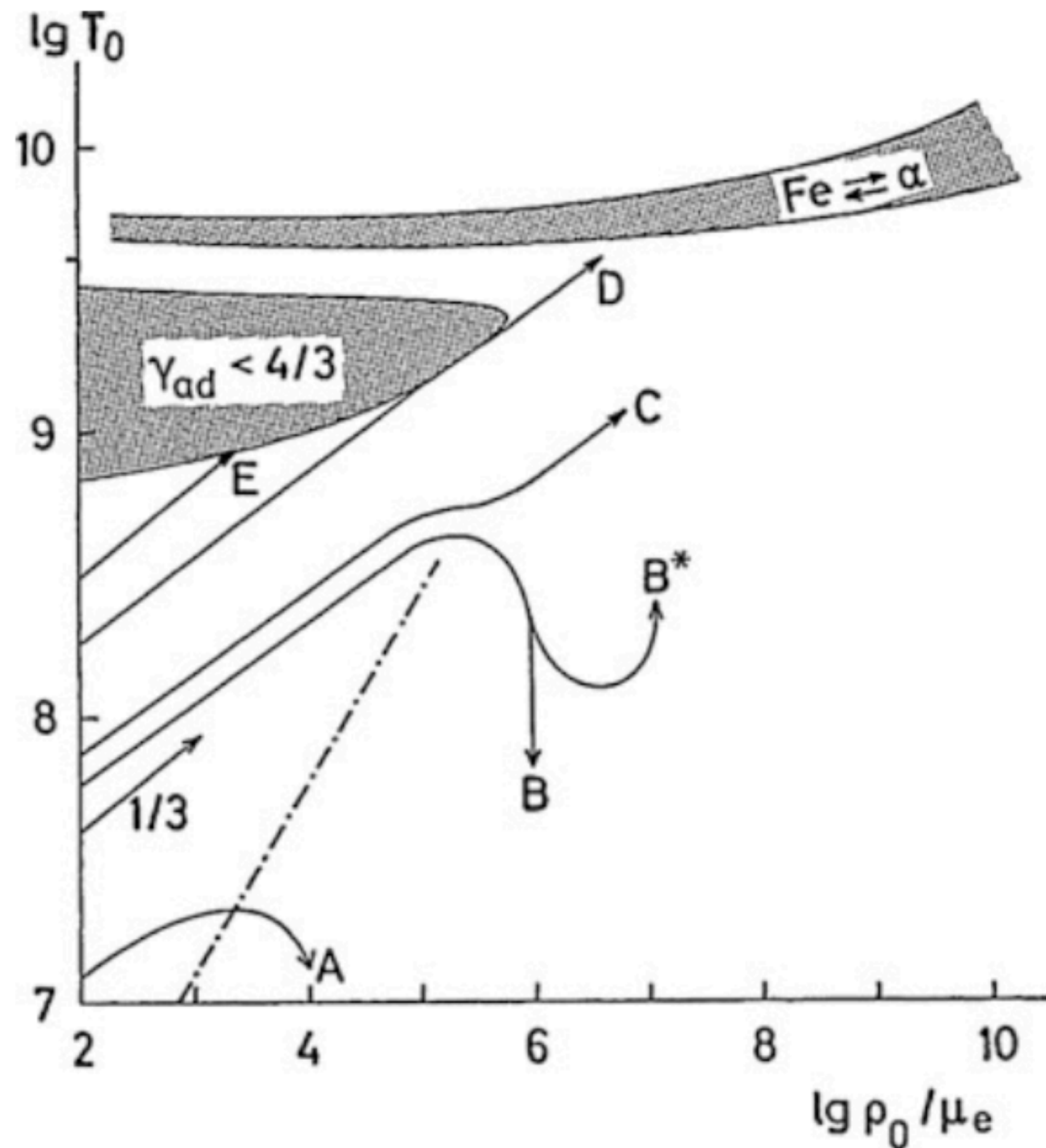
藤林 翔（東北大）

in collaboration with
S. Wanajo, A. T.-L. Lok,
C. Jockel, K. Kawaguchi, K. Ioka, and M. Shibata



Fate of massive stars

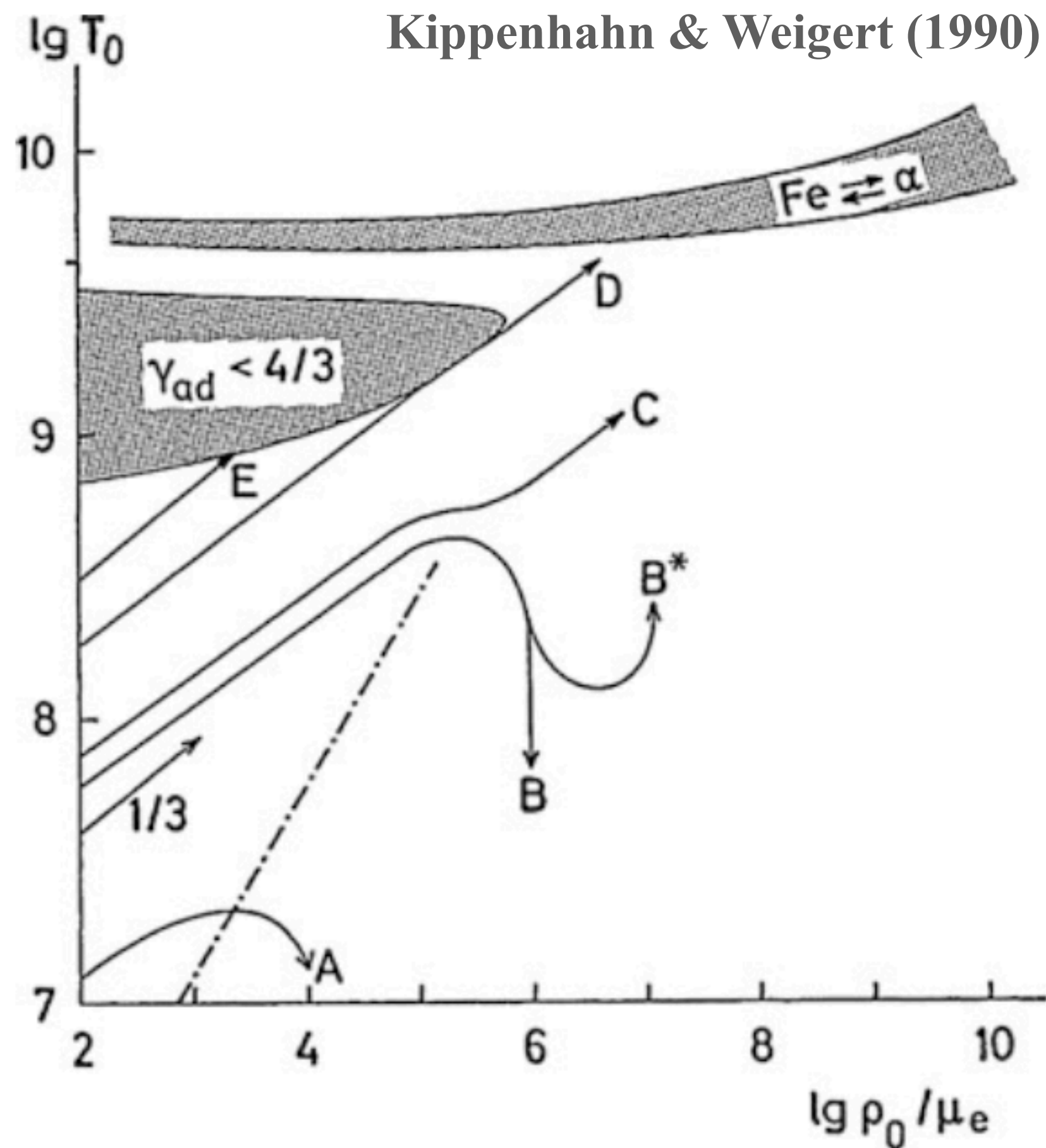
SMS



For single, non-rotating star...

- $M_{\text{ZAMS}} \lesssim (8 - 10)M_{\odot}$
→ Degeneracy pressure support before Fe-formation
- $M_{\text{ZAMS}} \gtrsim (8 - 10)M_{\odot}$
→ Fe-core formation → Gravitational collapse
- $M_{\text{ZAMS}} \gtrsim 130M_{\odot}$
→ e^-e^+ pair production → Gravitational collapse
- $M_{\text{ZAMS}} \gtrsim 10^4M_{\odot}$
→ General relativistic instability → Gravitational collapse

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Outline

- ✓ Collapse of rotating **supermassive** stars ($\sim 10^5 M_{\text{sun}}$, $\sim 10^3 M_{\text{sun}}$)
- ✓ Collapse of rotating massive star (System $\sim 10 M_{\text{sun}}$, ejecta $\sim M_{\text{sun}}$)

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Supermassive star

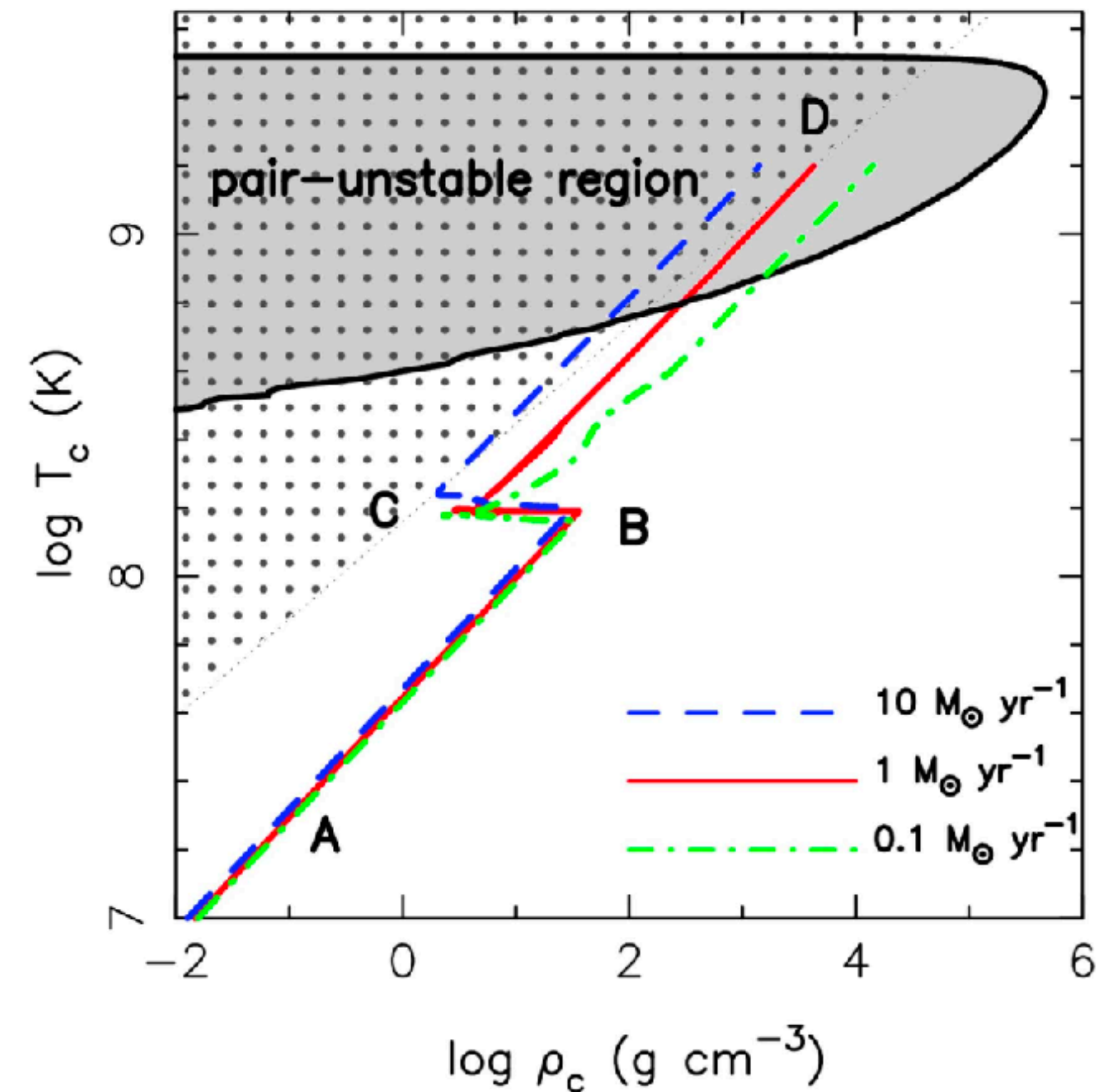
- Hypothetical very massive ($\gtrsim 10^4 M_\odot$) star
- High- s , P_{rad} -dominant ($\Gamma \approx 4/3$)
- Dies likely by GR instability

$$\rho_{\text{crit}} \approx 1.994 \times 10^{18} \left(\frac{0.5}{\mu}\right)^3 \left(\frac{M_\odot}{M}\right)^{7/2} \text{ g cm}^{-3}$$

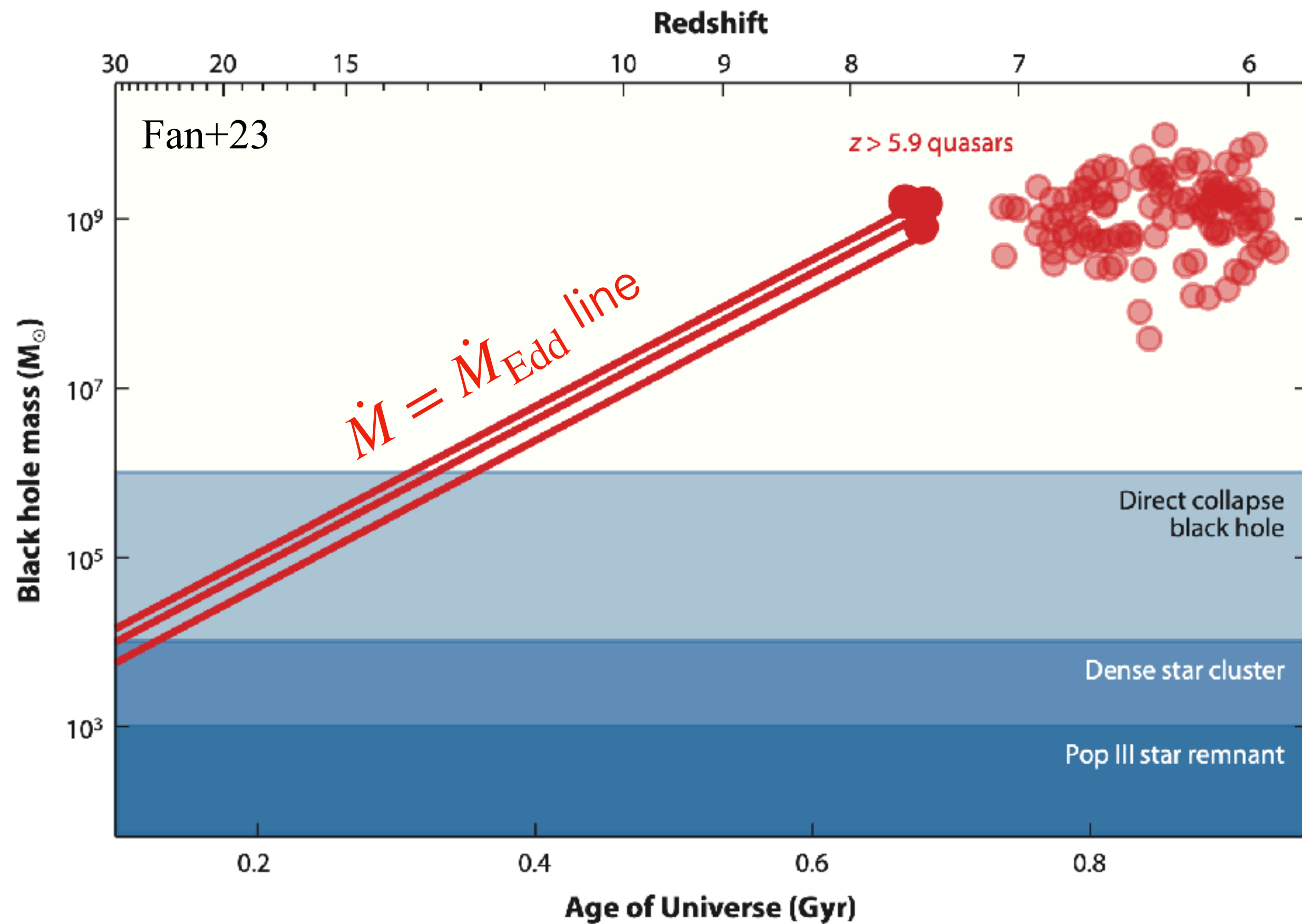
$$\approx 4 \text{ g/cm}^3 (M/10^5 M_\odot)^{-3.5} (\mu/0.59)^{-3}$$

(Shapiro-Teukolsky 83, Fuller+86)

Umeda+16



SMBHs in early universe



Distant (high- z) SMBH with $M \sim 10^9 M_{\odot}$

How are they formed?

Basic idea: $\dot{M} \lesssim \dot{M}_{\text{Edd}} \propto M$.

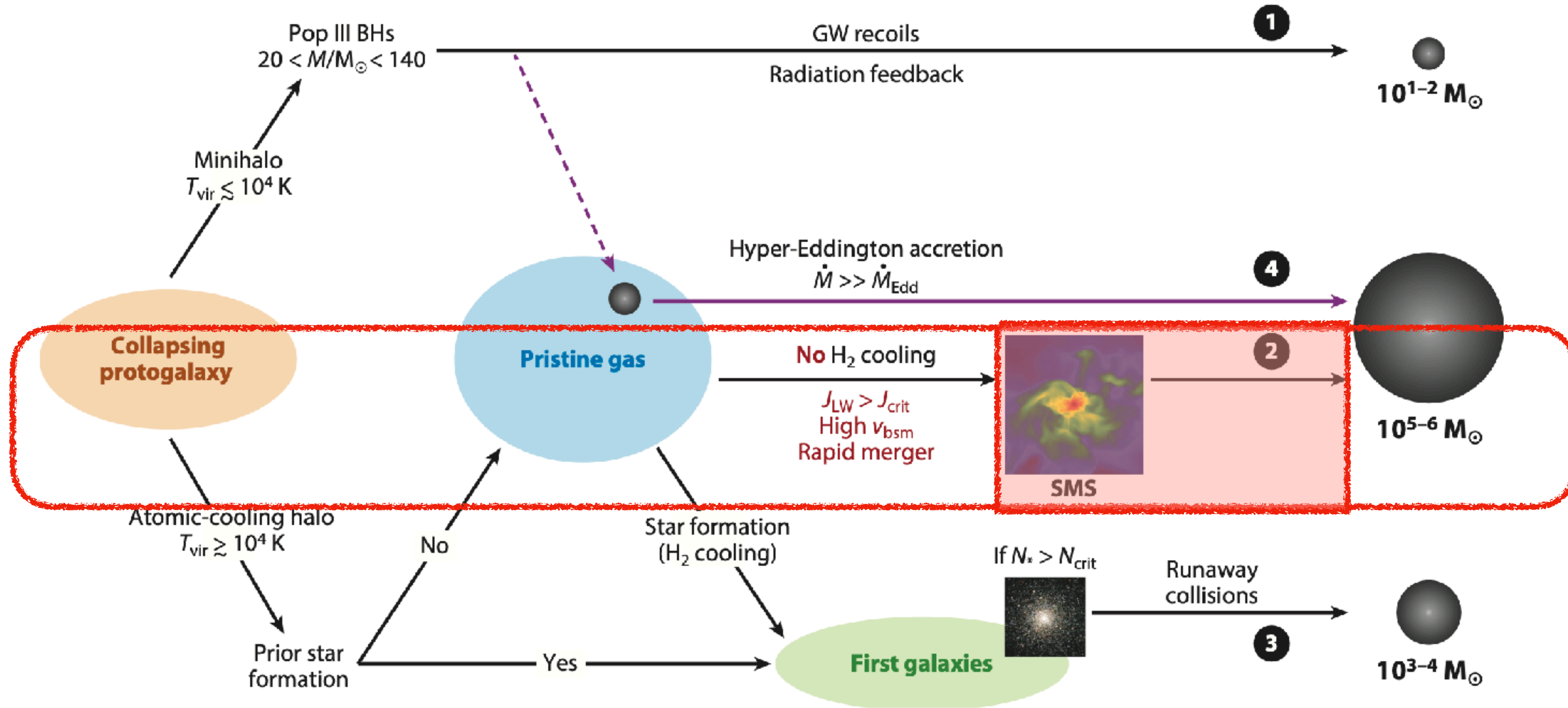
(Super/hyper-Eddington accretion may be possible)

Keeping a high Eddington ratio $\dot{M}/\dot{M}_{\text{Edd}}$

from $10^2 M_{\odot}$ to $10^9 M_{\odot}$ may not be easy

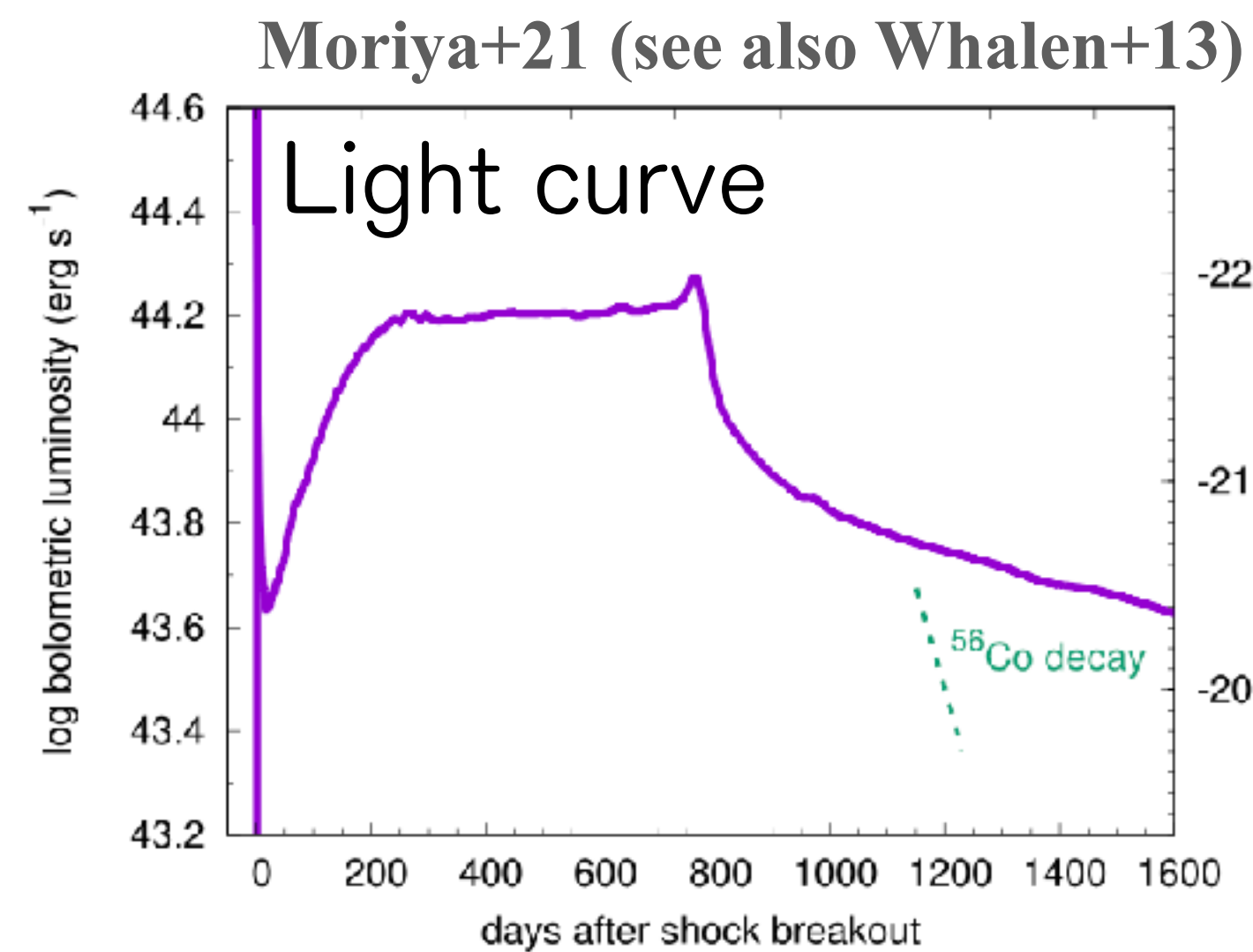
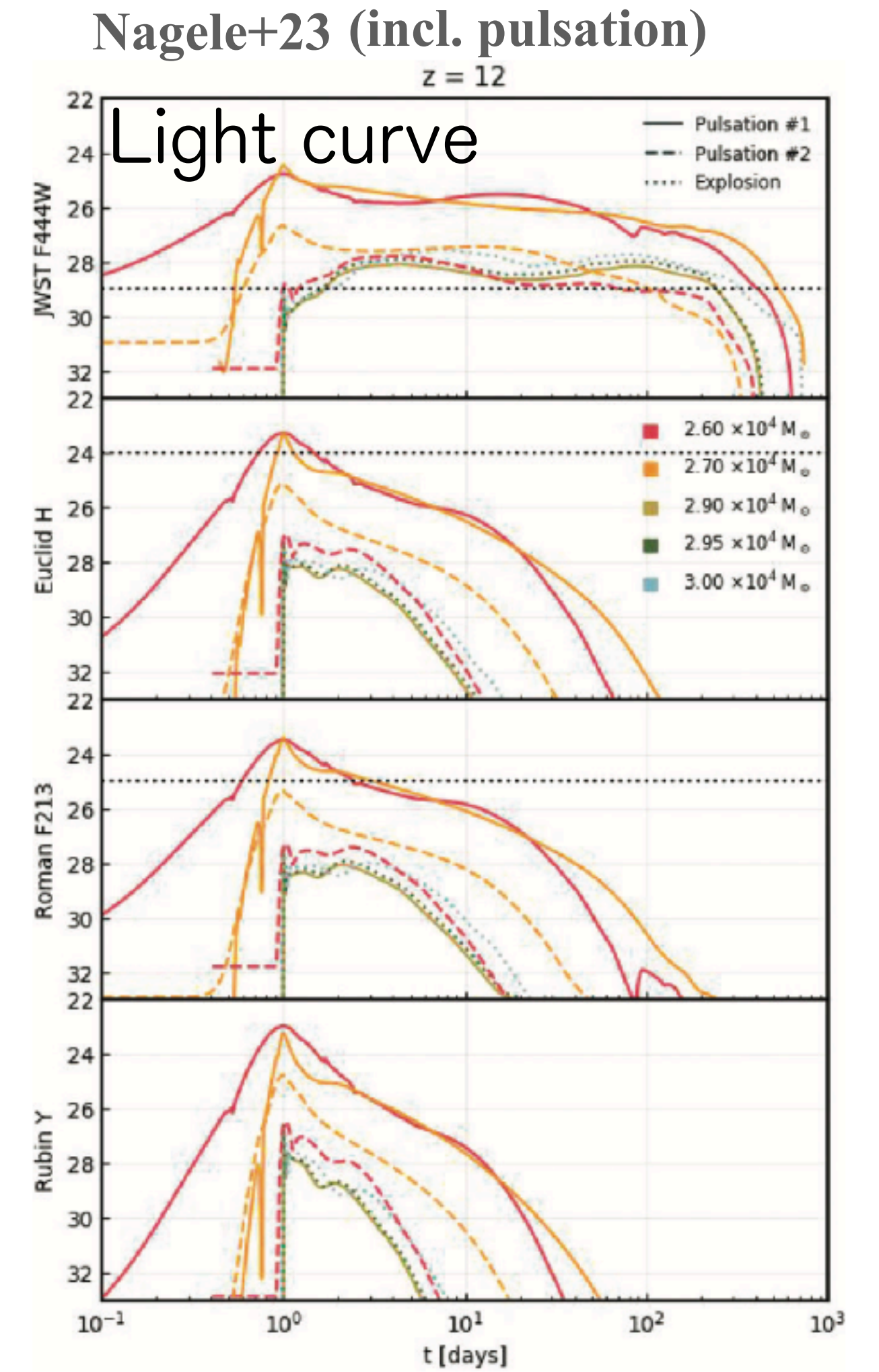
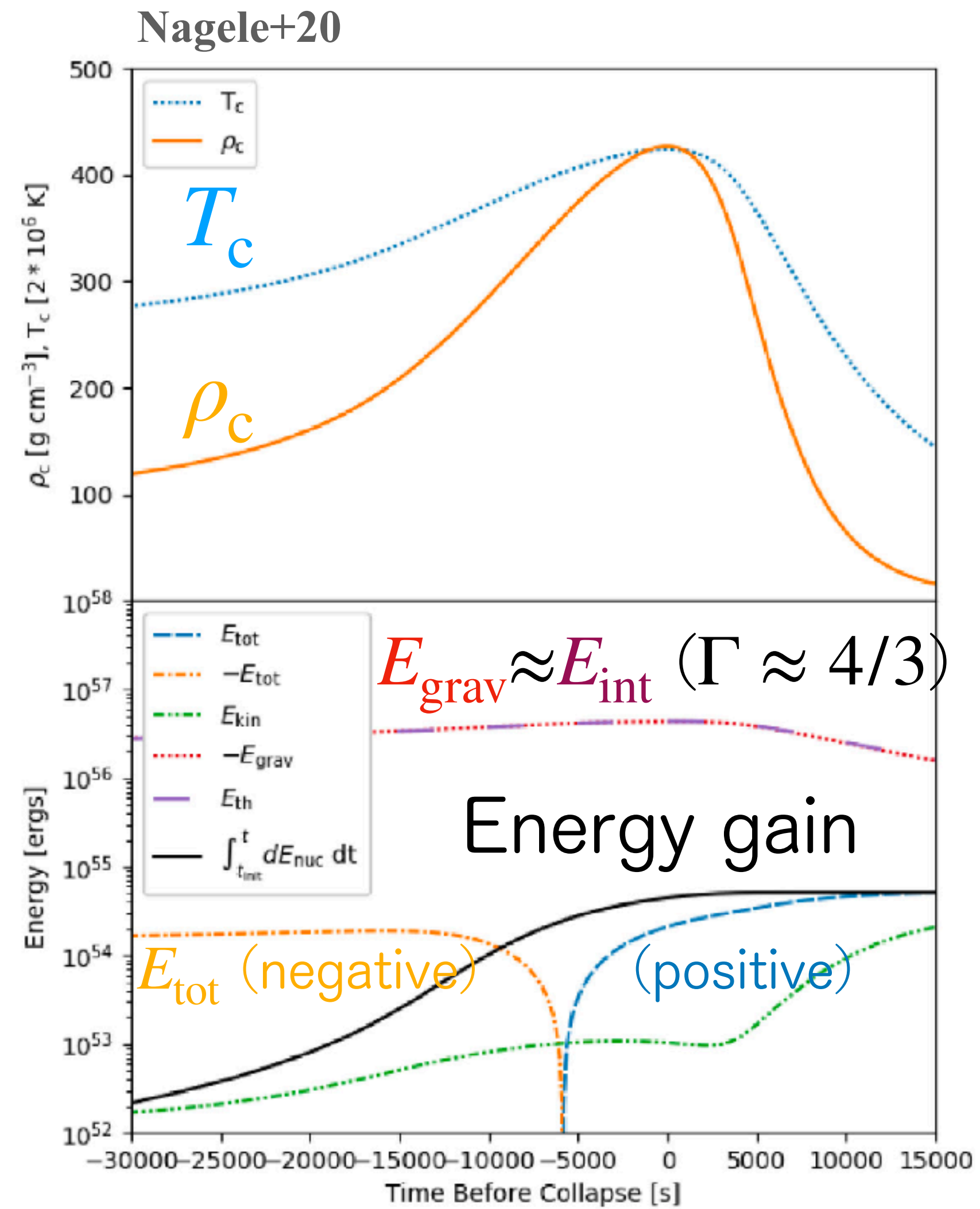
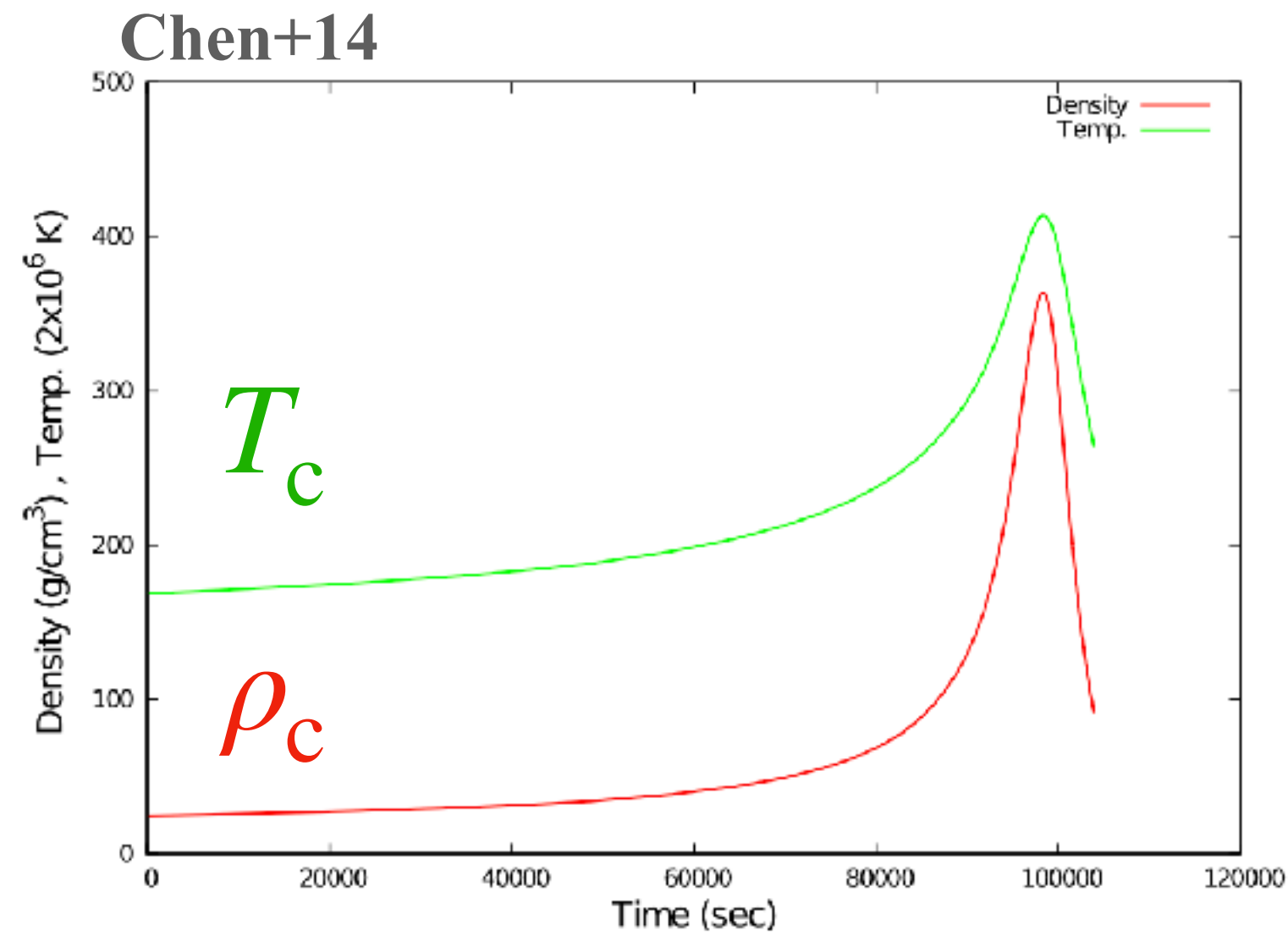
Initial high mass of BH may make it easier!

Direct collapse scenario for SMBHs in early universe



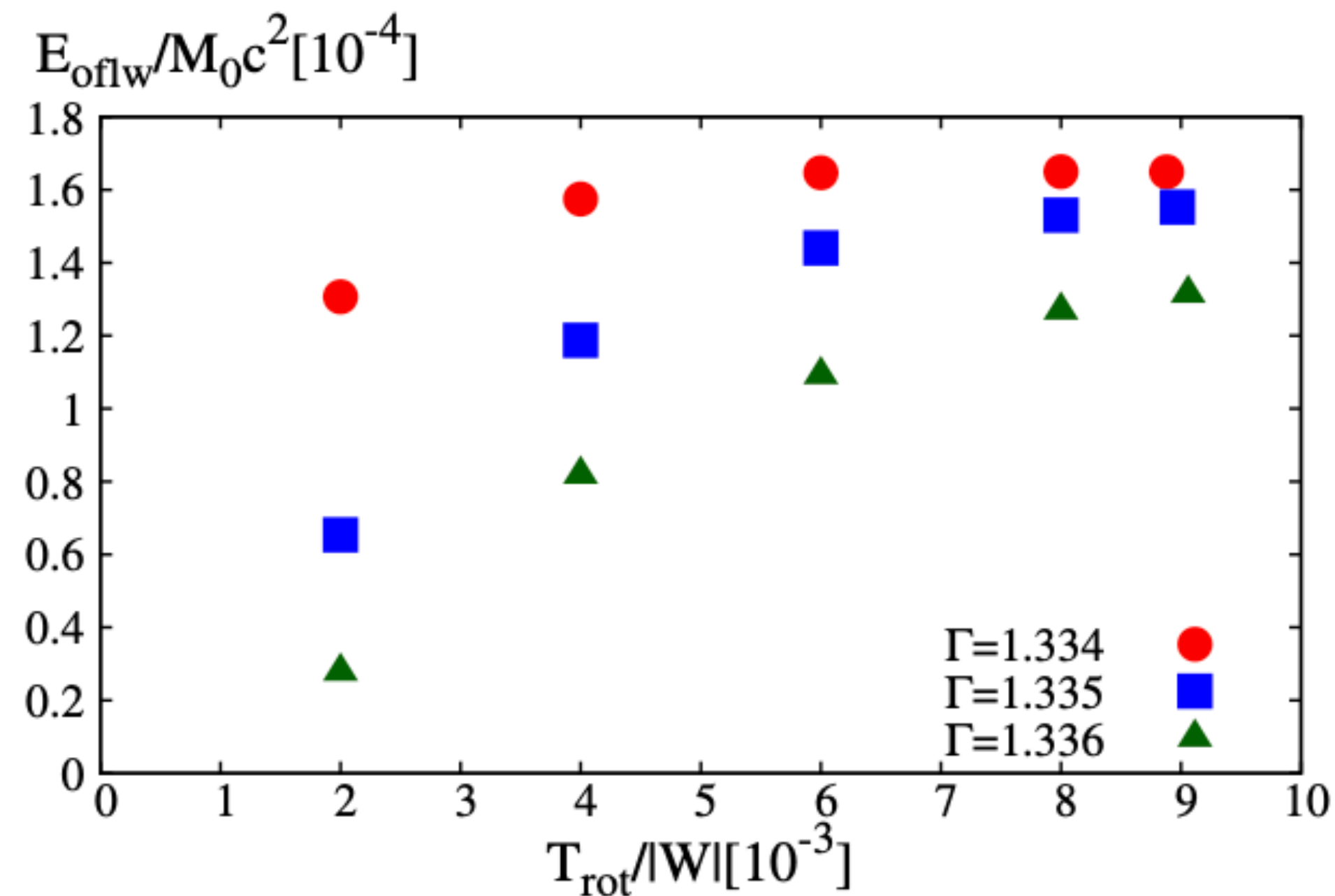
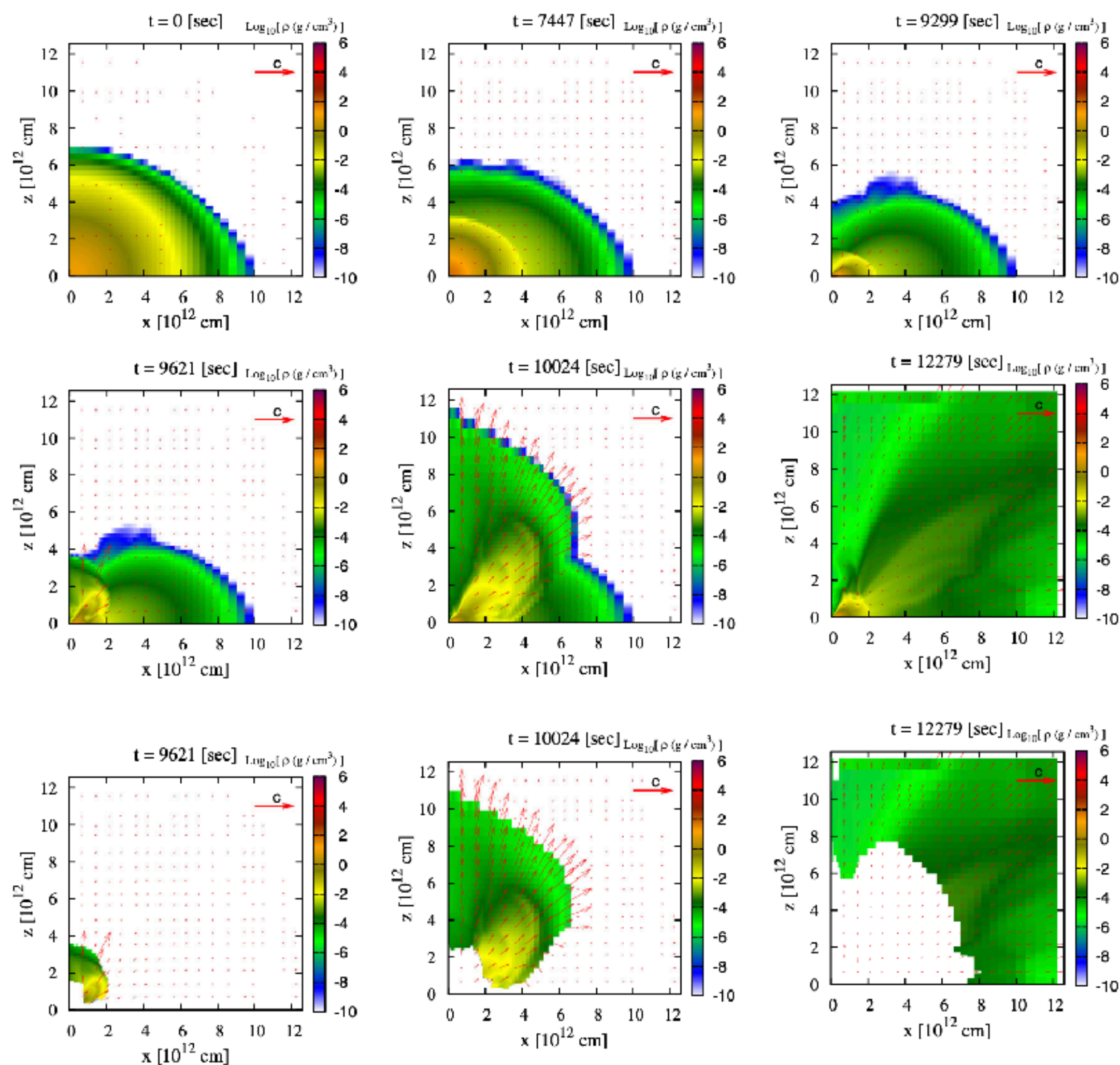
Inayoshi+20, see also Rees 1978

Explosion of supermassive stars: thermonuclear case



Effects of rotation: bounce-induced explosion

Uchida+17



Let us revisit this scenario as the first step toward the understanding collapse of very massive stars

Method: Numerical setup

General relativistic gravity

Hydrodynamics with nuclear reaction

$H \rightarrow He \rightarrow C$ (Only forward reaction)

CNO cycle triple- α

Equation of state

Composite of ions(H, He, C), photons, electrons and positrons

Neutrino radiation

Only for neutrinos emitted by CNO cycle $\sim 8\%$ of heating rate

Method: Initial supermassive star models

Marginally stable SMS with rotation.

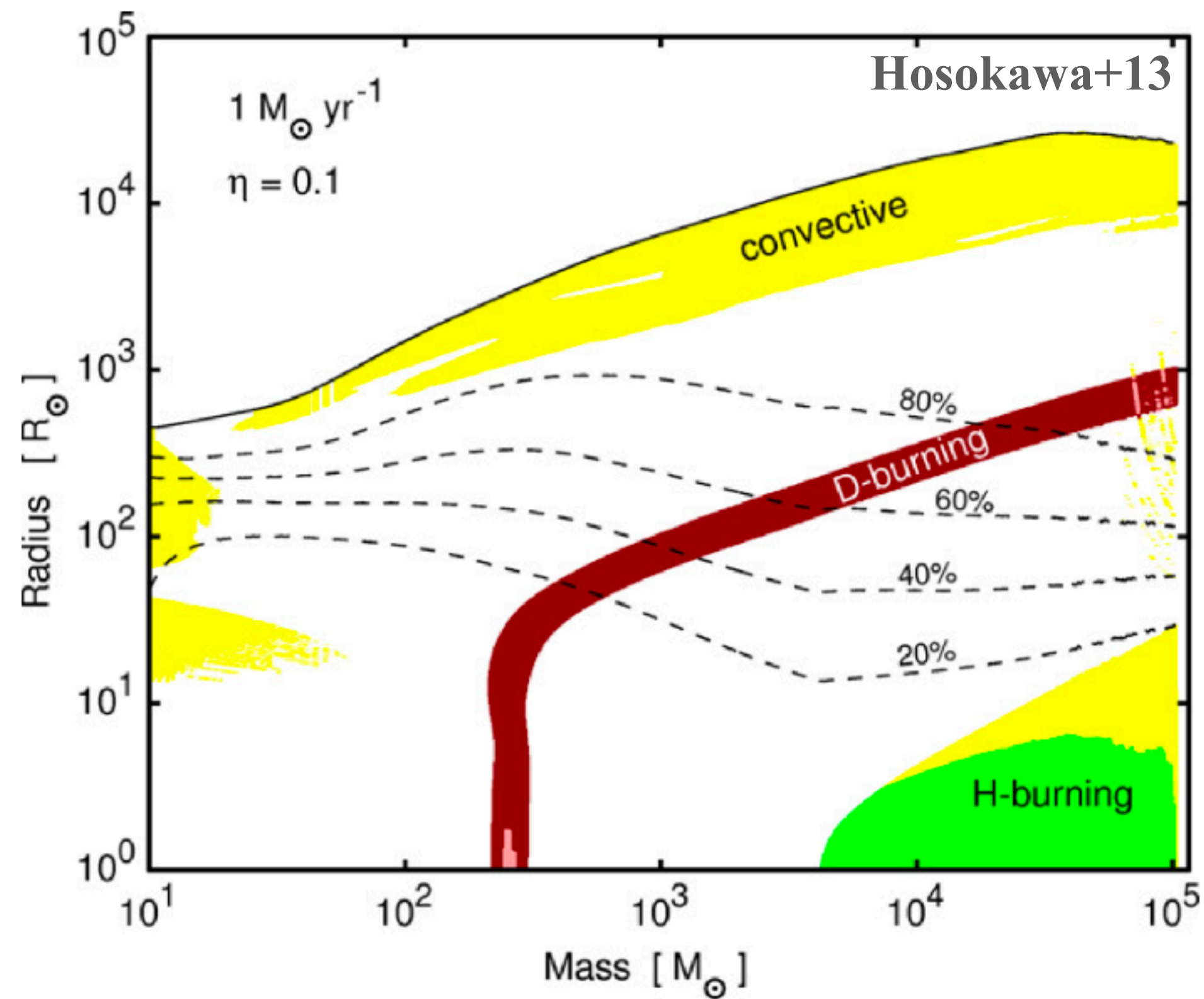
model	$M_0 (M_\odot)$	$R_{e0} \text{ (cm)}$	$T_{\text{kin}}/ W $	$\alpha_{c,0}$	$\gamma_{c,0} - 4/3$	\hat{A}	s/k_B
H1	2.1×10^5	1.7×10^{13}	0.002	0.992	0.0026	∞	450
H2	3.2×10^5	2.3×10^{13}	0.004	0.990	0.0021	∞	550
H3	4.3×10^5	2.7×10^{13}	0.006	0.988	0.0018	∞	630
H4	6.9×10^5	4.4×10^{13}	0.009	0.985	0.0014	∞	800
Hdif1	9.2×10^5	5.0×10^{13}	0.011	0.983	0.0012	2	920
Hdif2	1.1×10^6	5.3×10^{13}	0.013	0.981	0.0012	1.5	1000
Hdif3	1.9×10^6	7.4×10^{13}	0.018	0.976	0.0009	1.0	1300
He1	5.0×10^4	4.3×10^{12}	0.002	0.992	0.0023	∞	210
He2	7.1×10^4	5.1×10^{12}	0.004	0.990	0.0019	∞	250
He3	9.6×10^4	6.1×10^{12}	0.006	0.988	0.0016	∞	300
He4	1.6×10^5	1.0×10^{13}	0.009	0.985	0.0013	∞	380

Primordial composition
 $X(\text{H})=0.25$, $X(\text{He})=0.75$

Purely He star
 $X(\text{He})=1$

Method: Initial supermassive star models

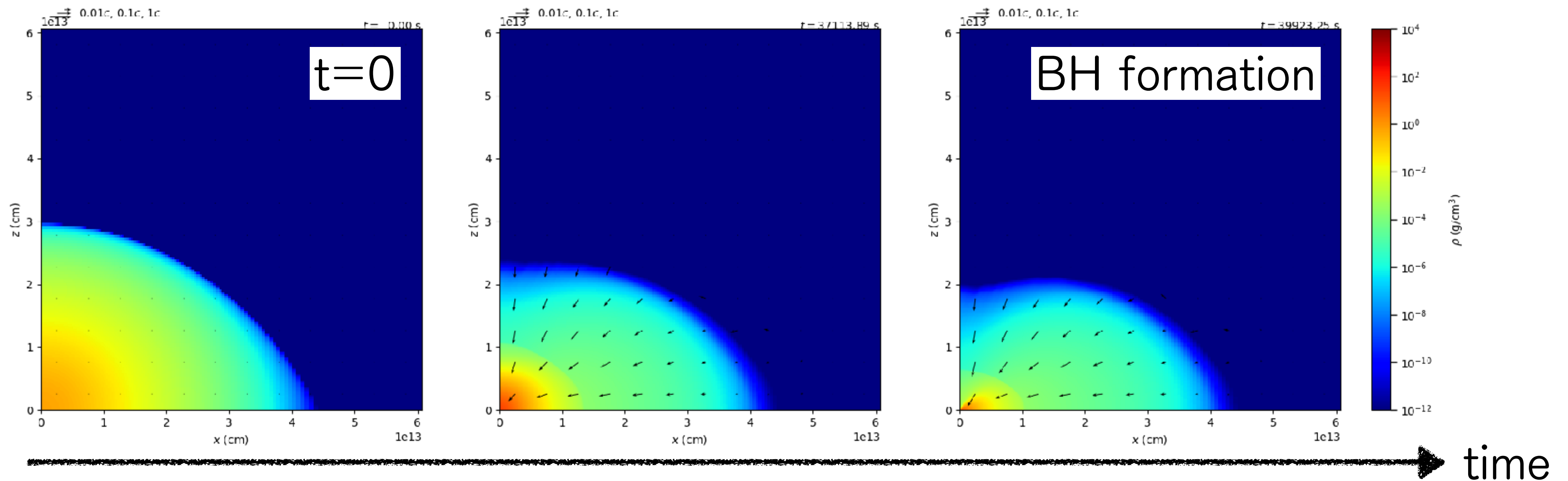
Caution: here are only the isentropic “core” of SMS



Realistic SMS may have inflated envelope

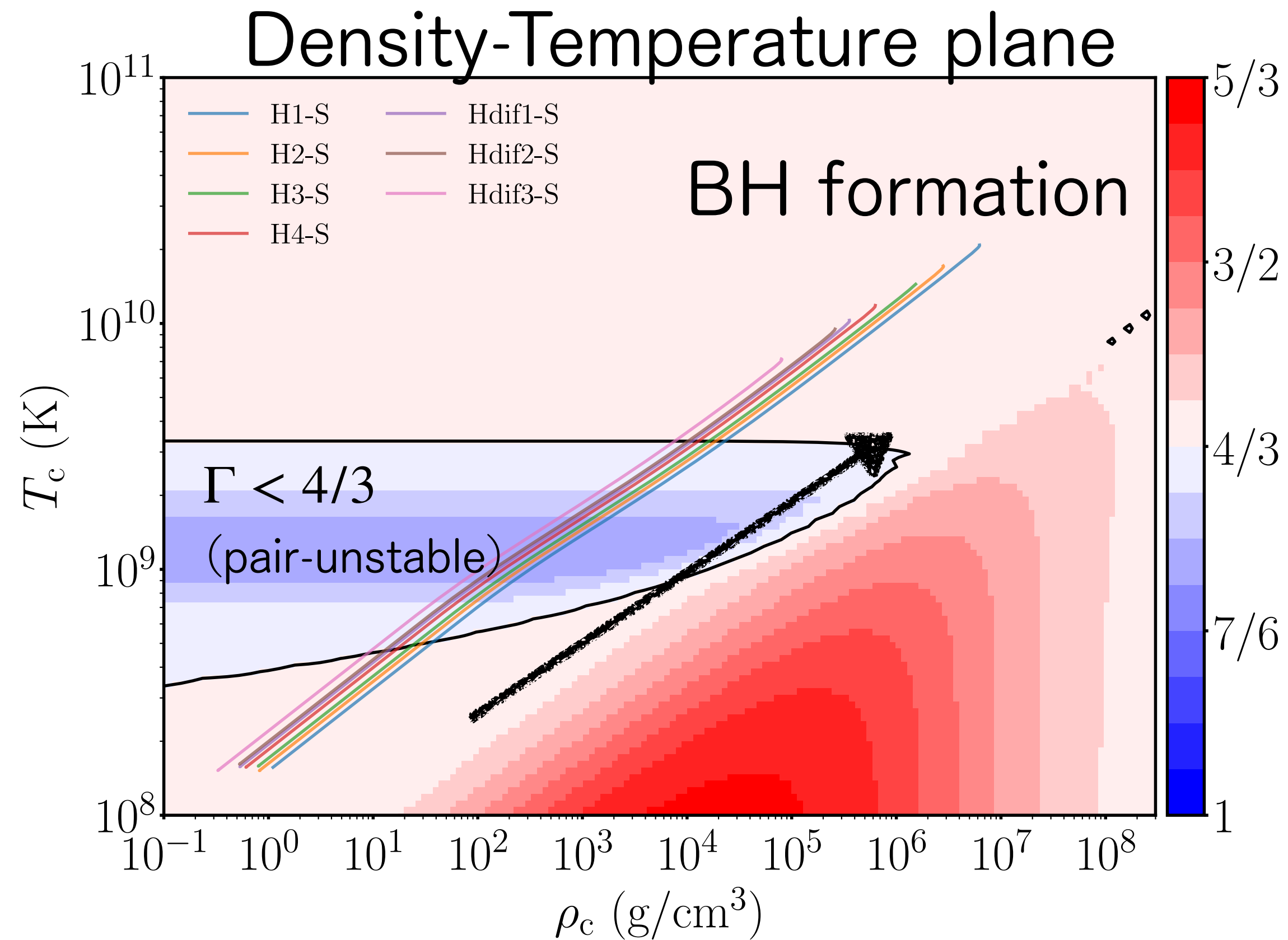
Result: Outline of evolution

Primordial composition, mass-shedding case



Collapsing motion is \sim coherent (characteristic of GR instability)

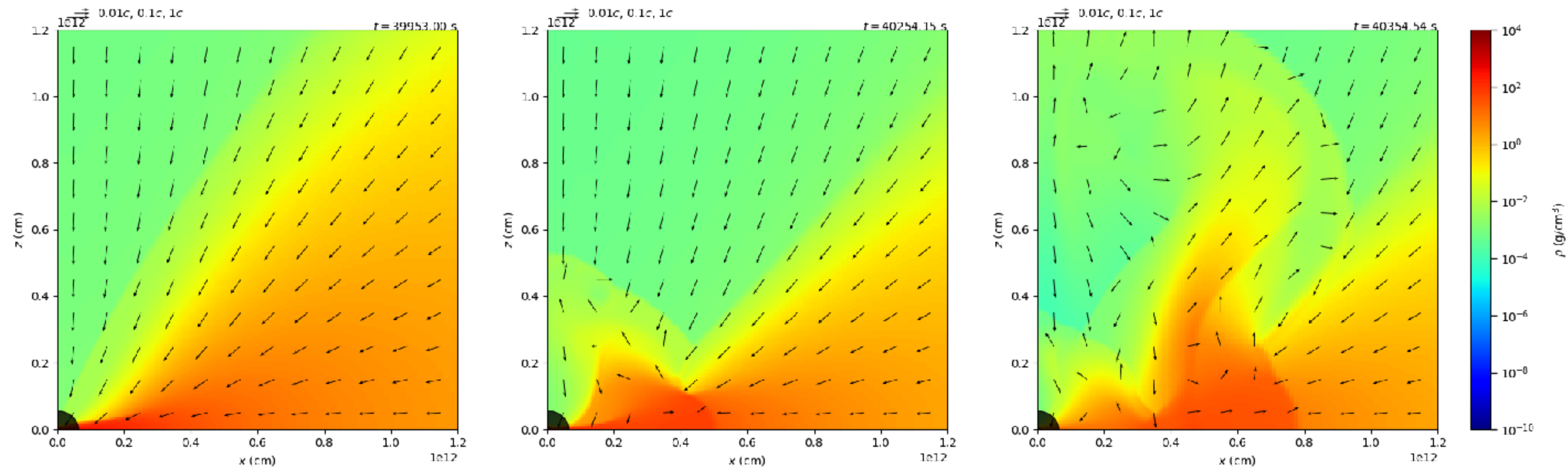
Result: Outline of evolution



Collapse proceeds
outside pair-unstable region.
→ GR instability

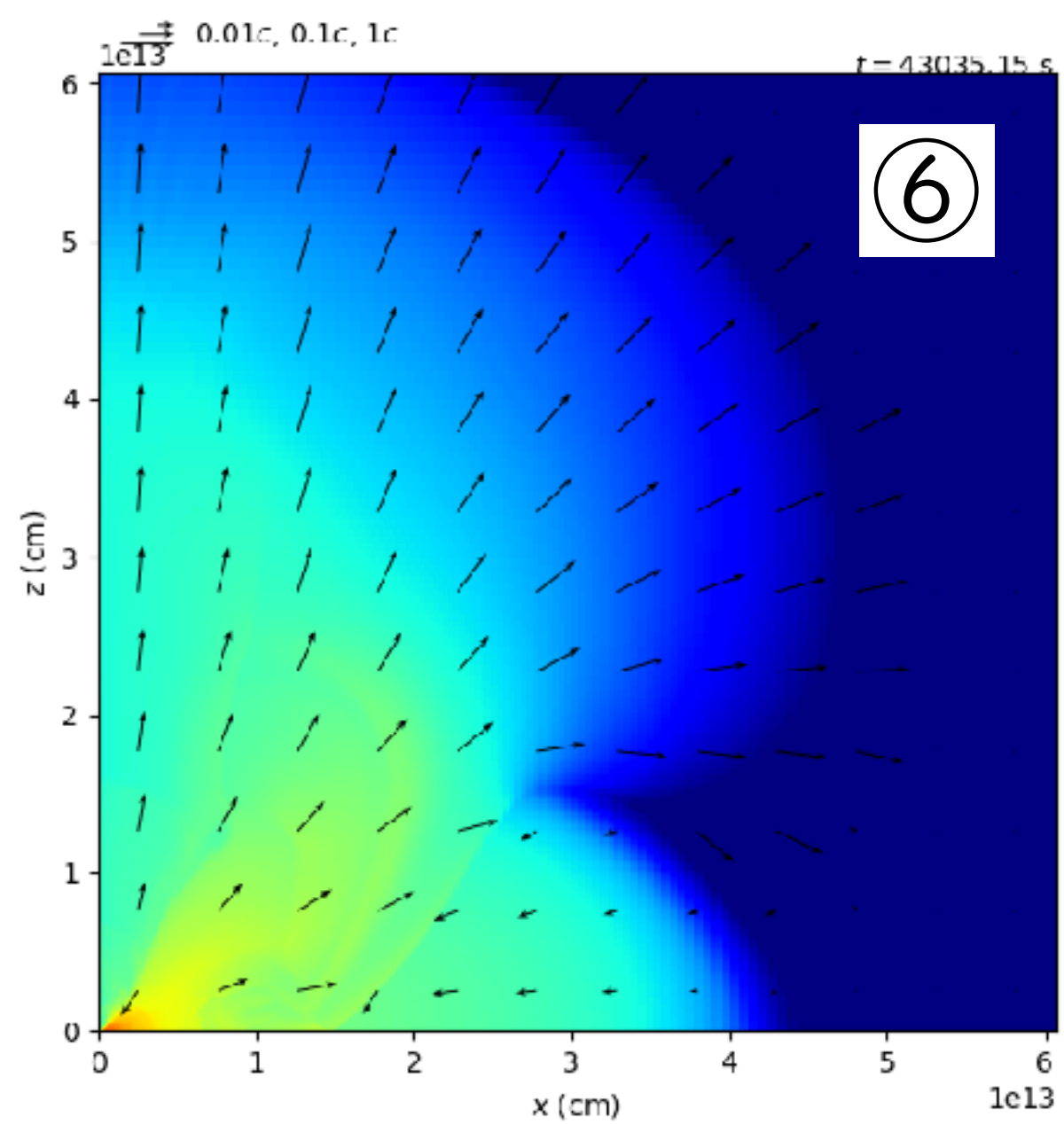
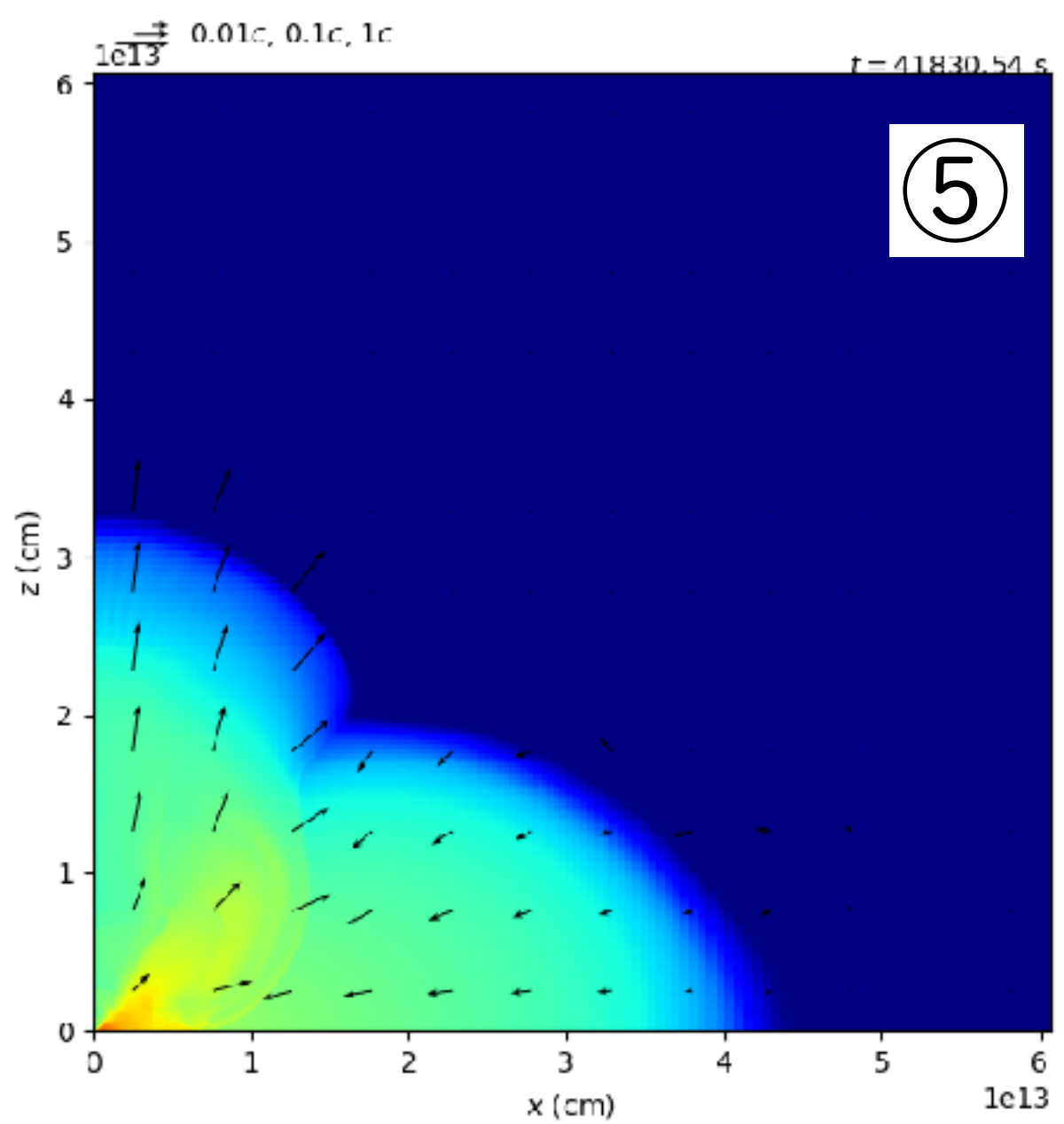
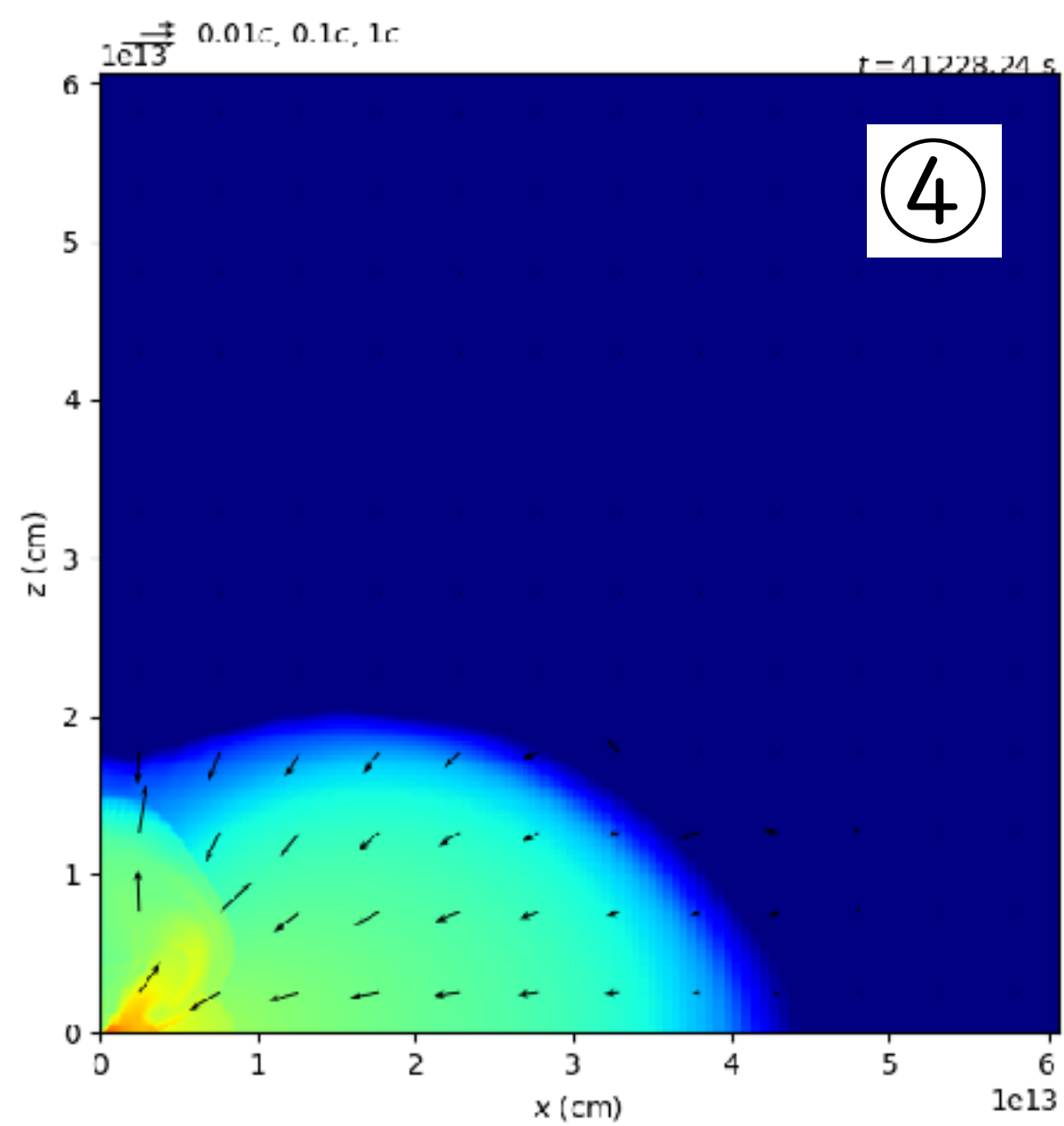
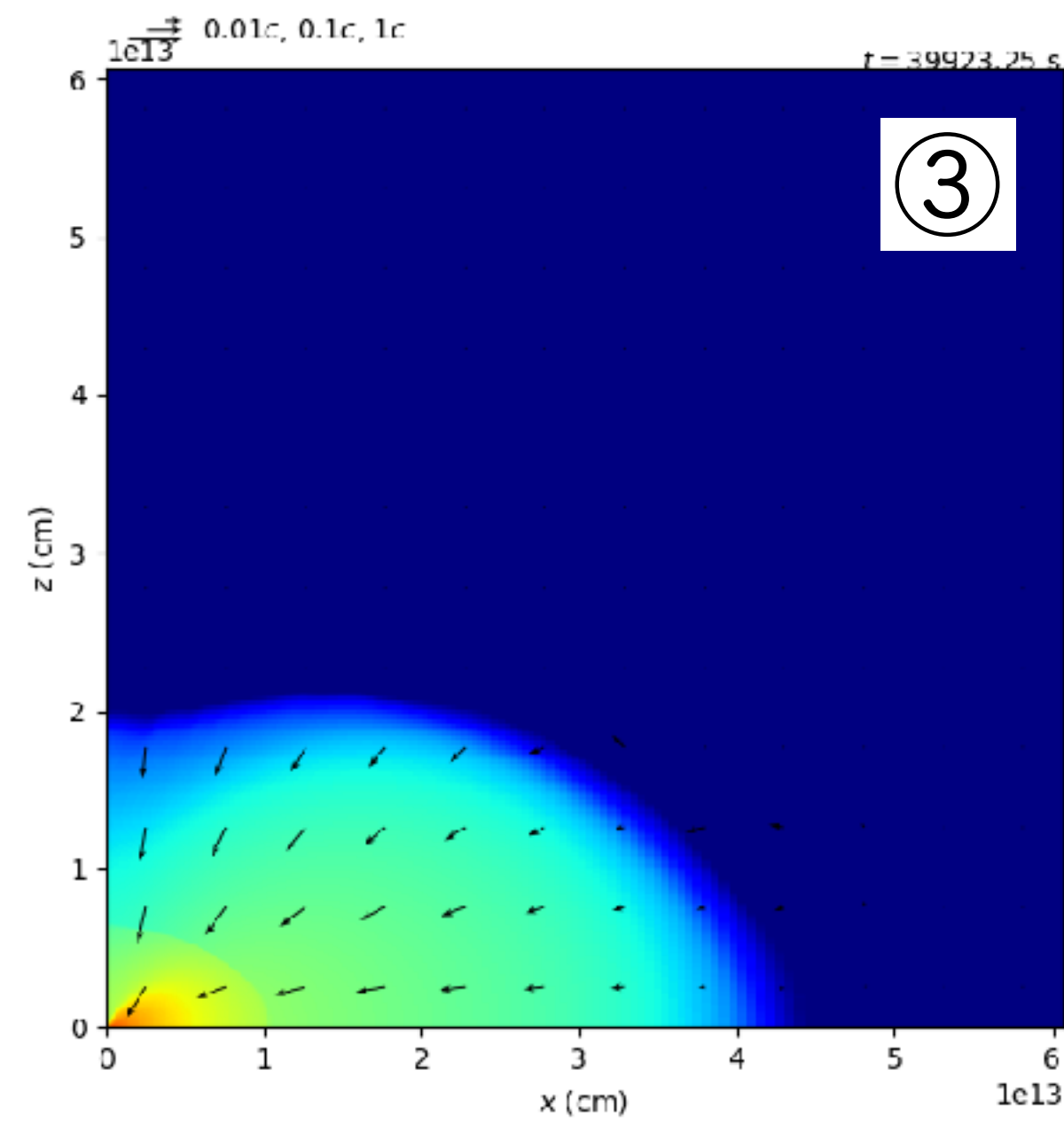
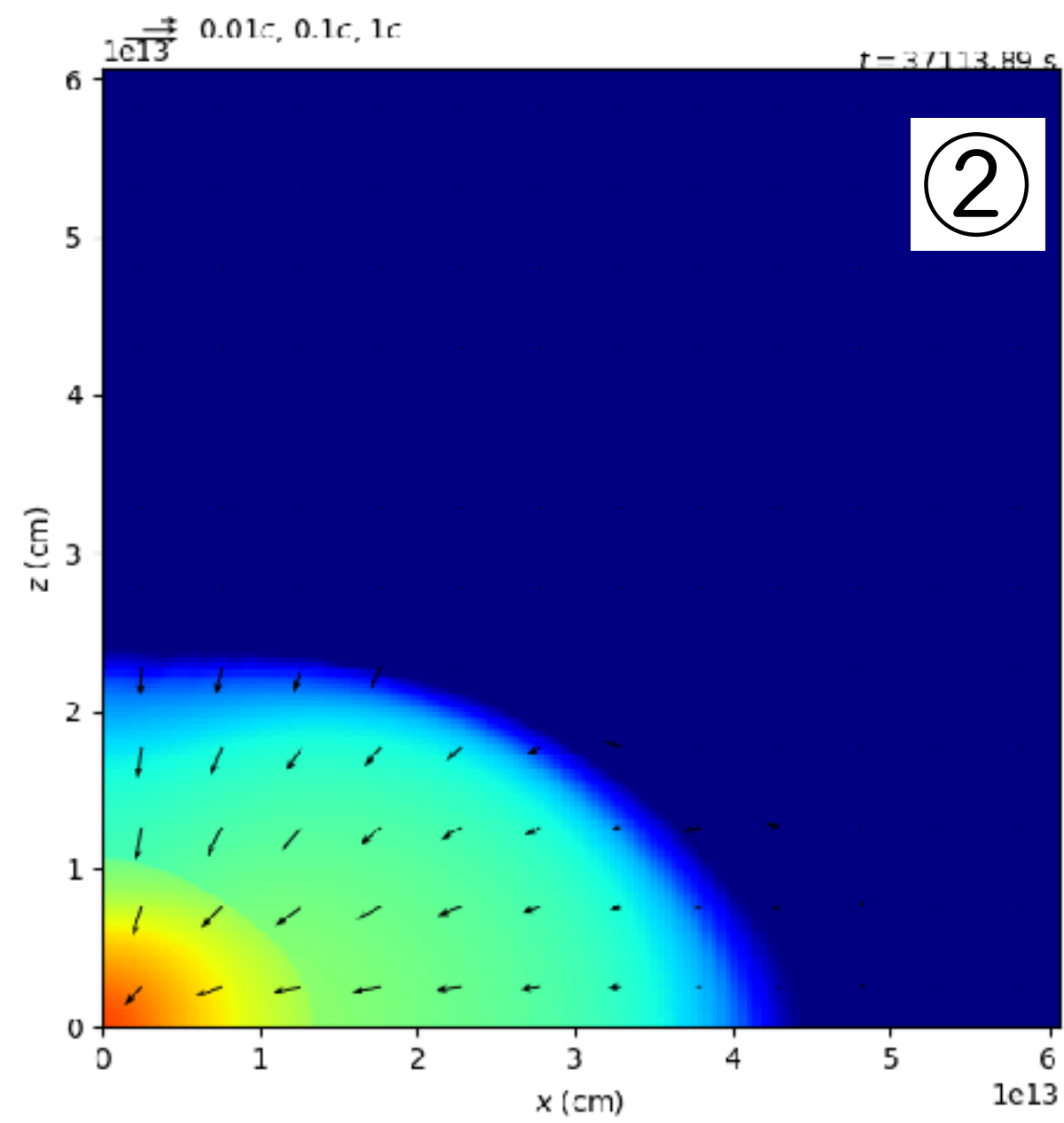
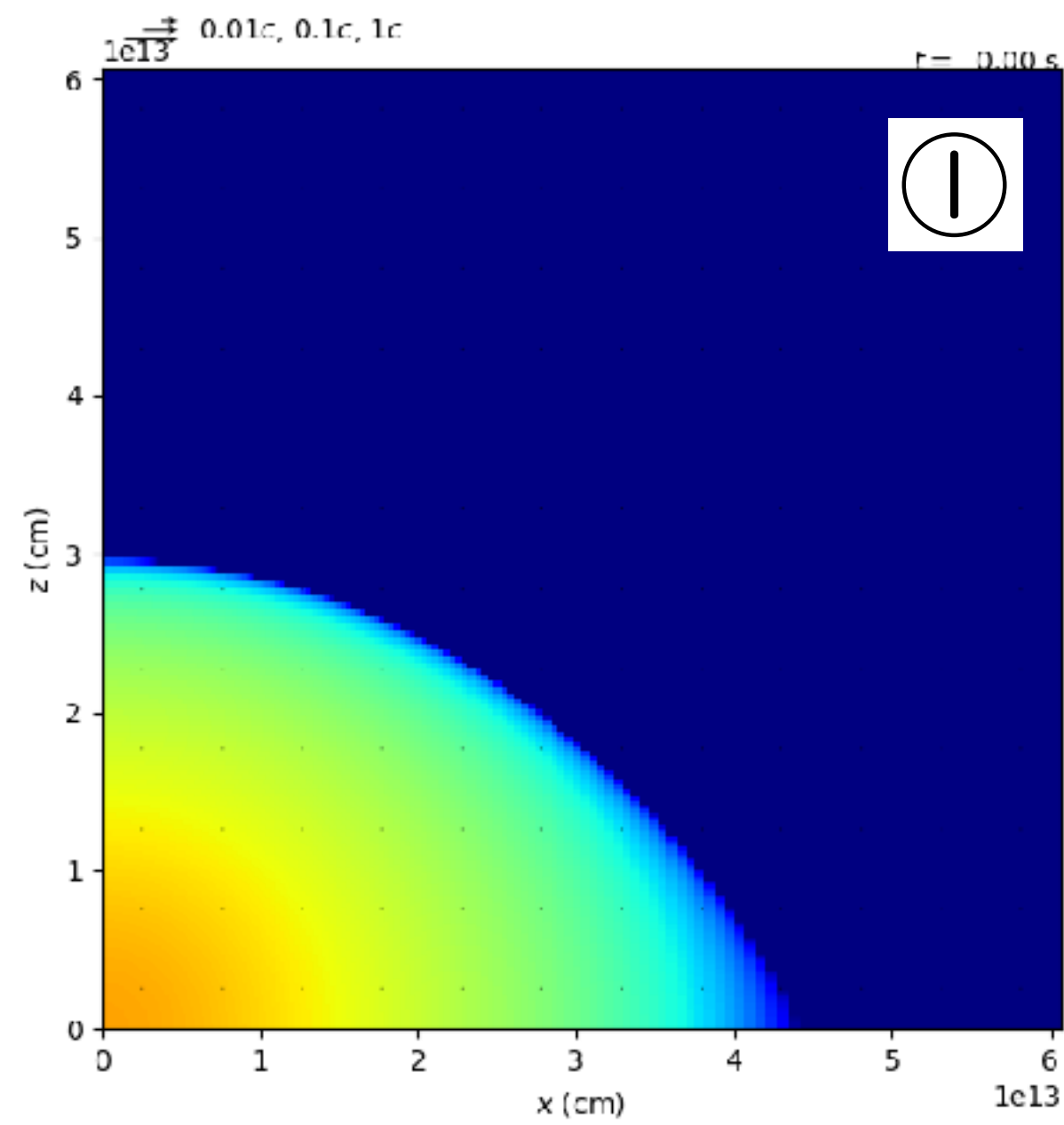
Result: Bounce-shock-induced ejecta

Density snapshots around torus formation time.

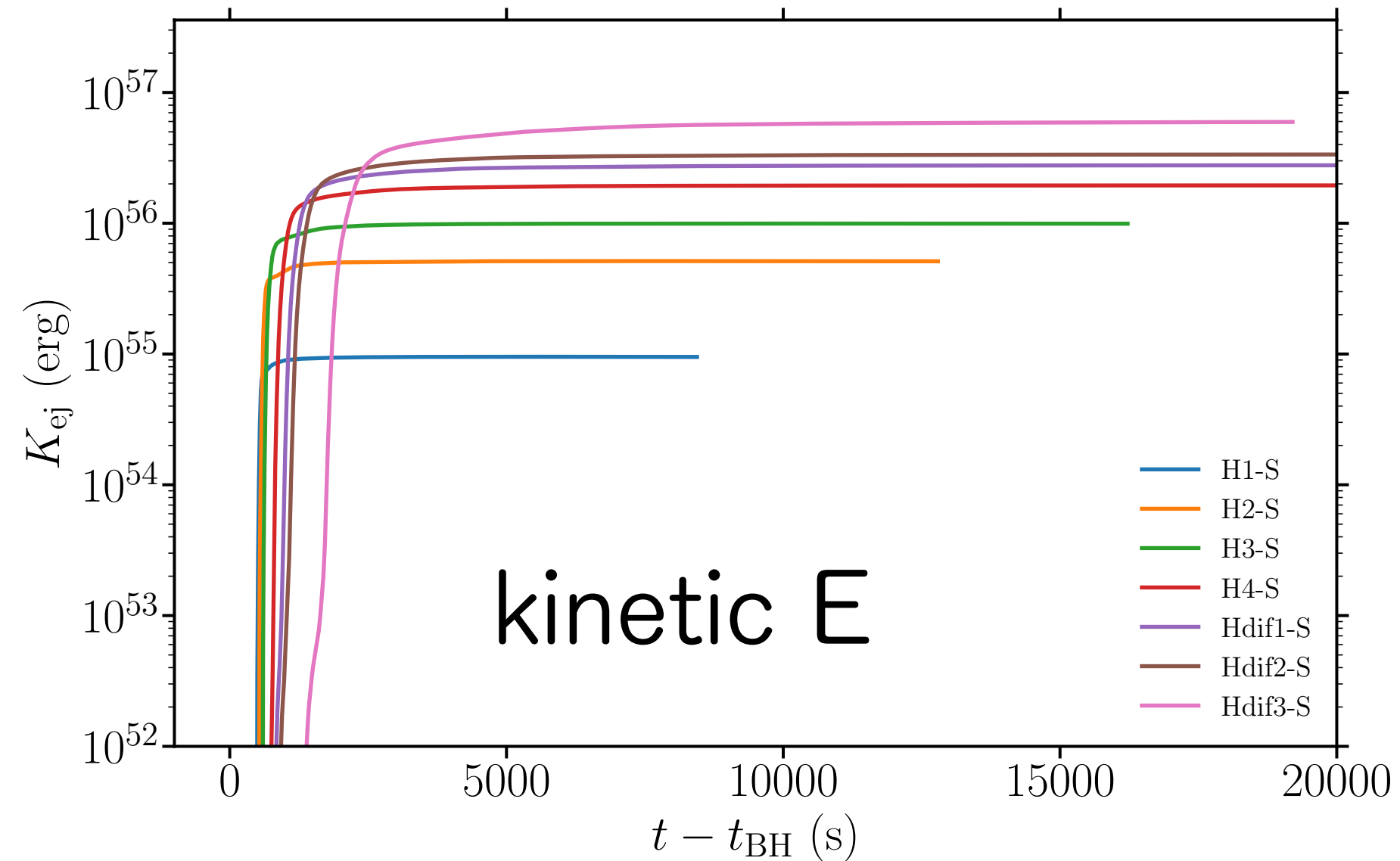
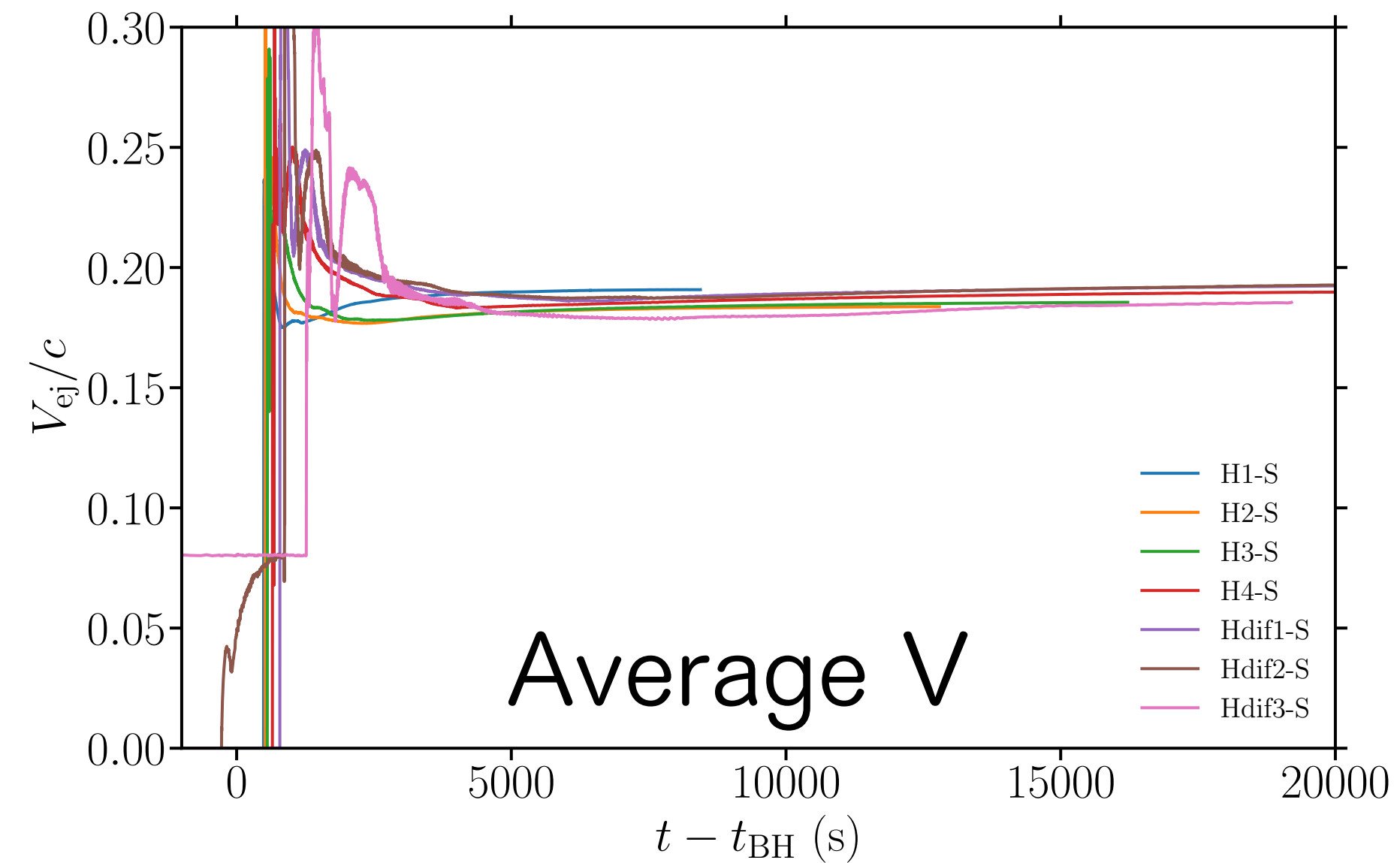
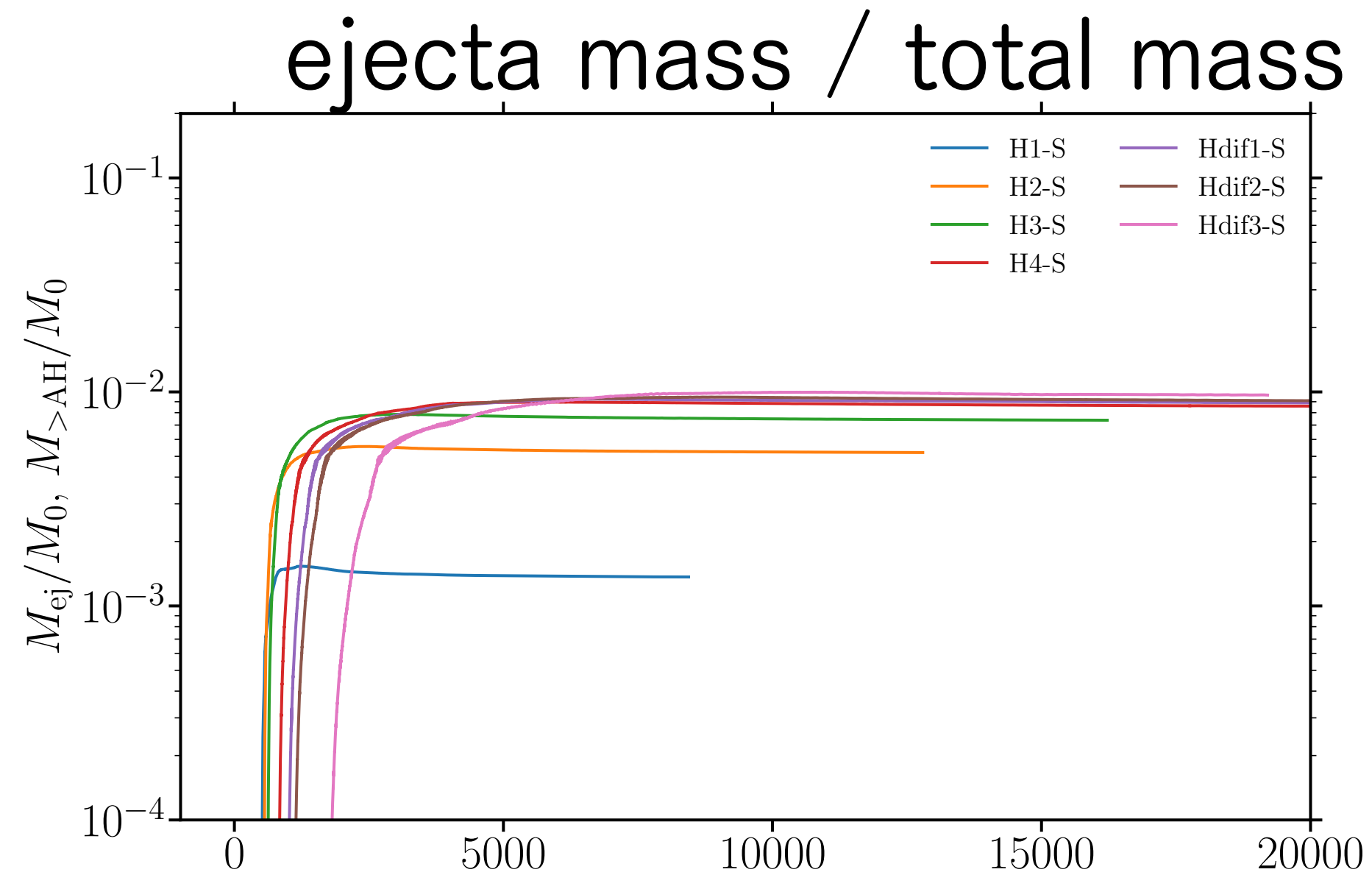


time

Sudden formation of centrifugally supported torus induces its bounce

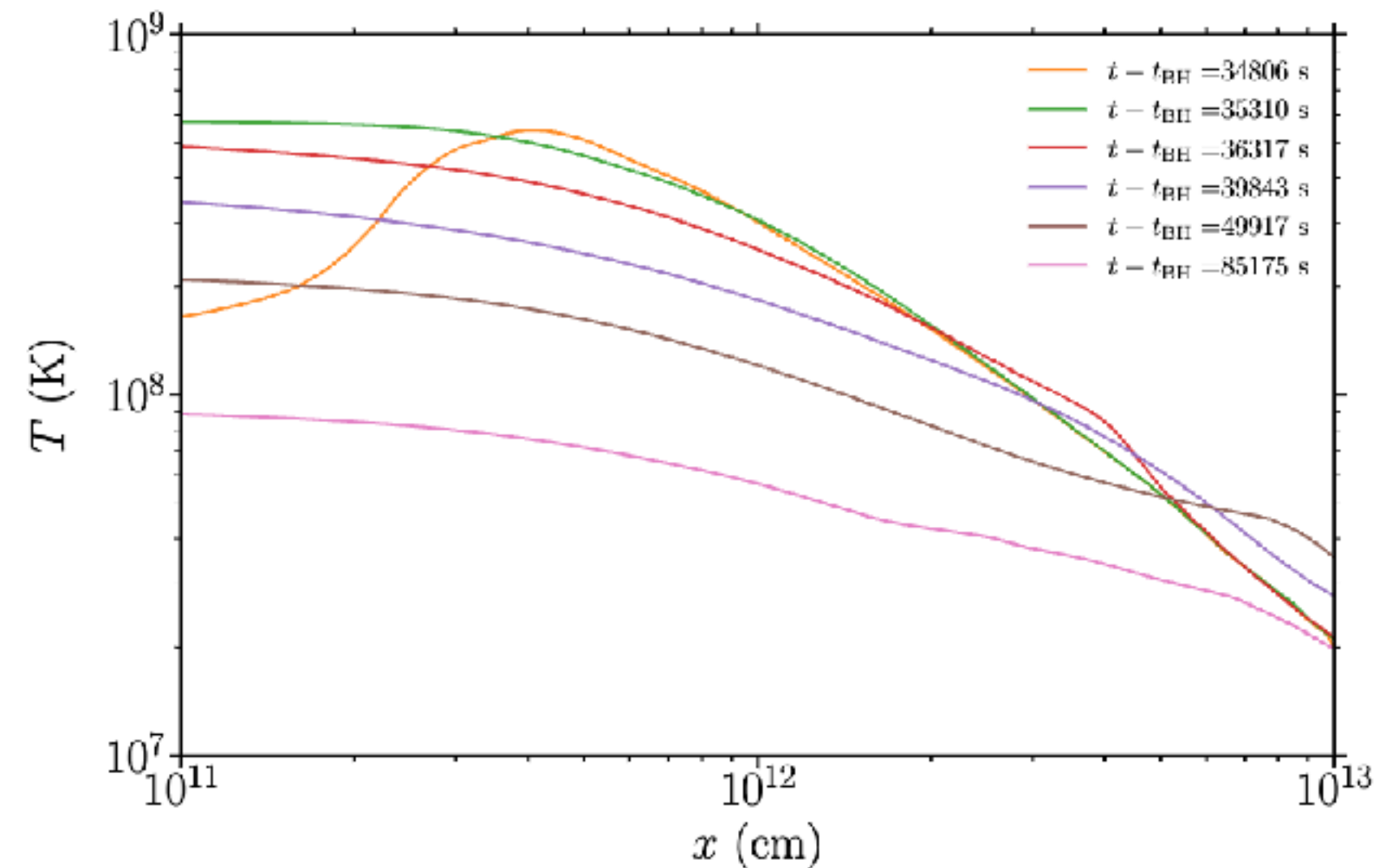
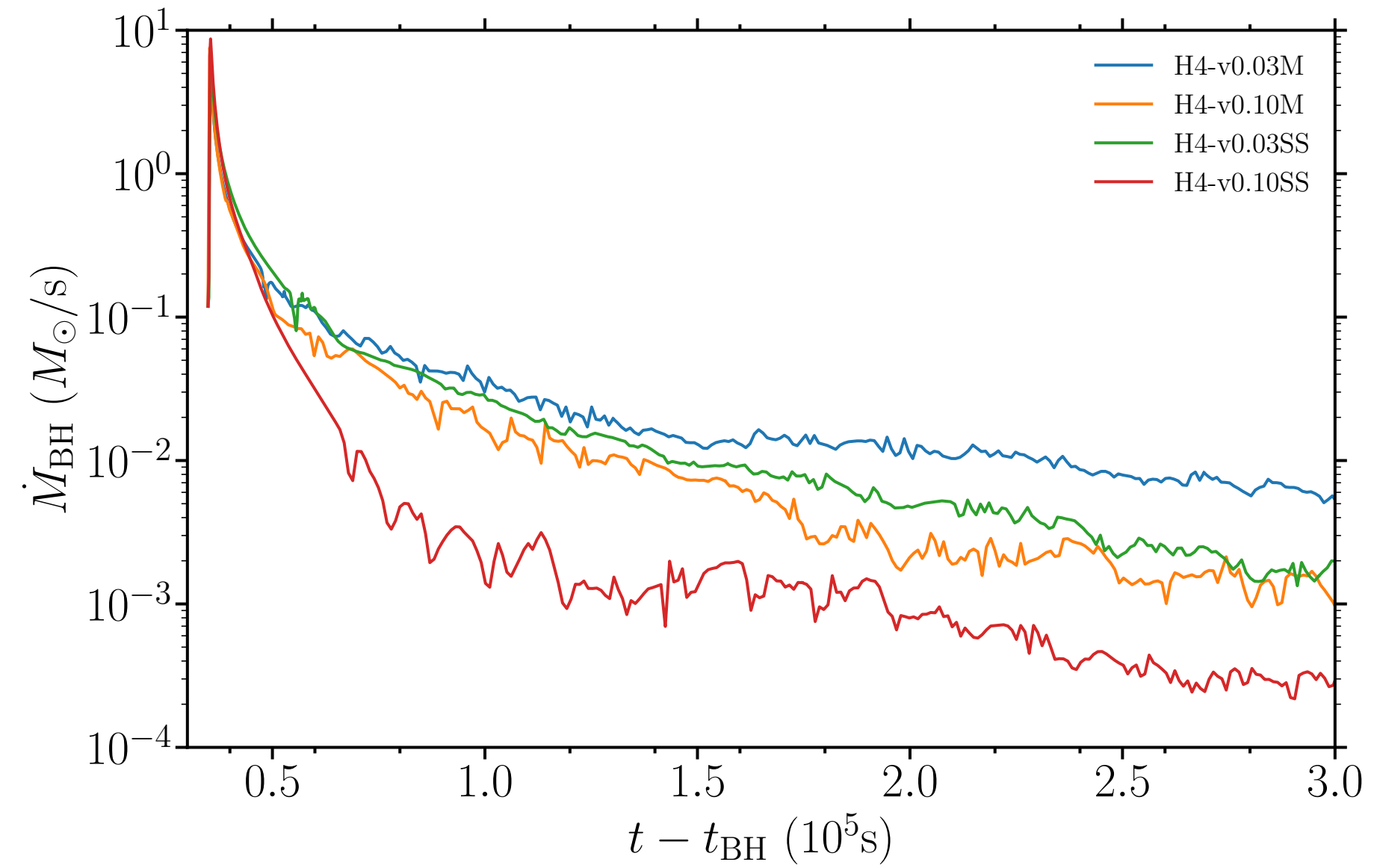
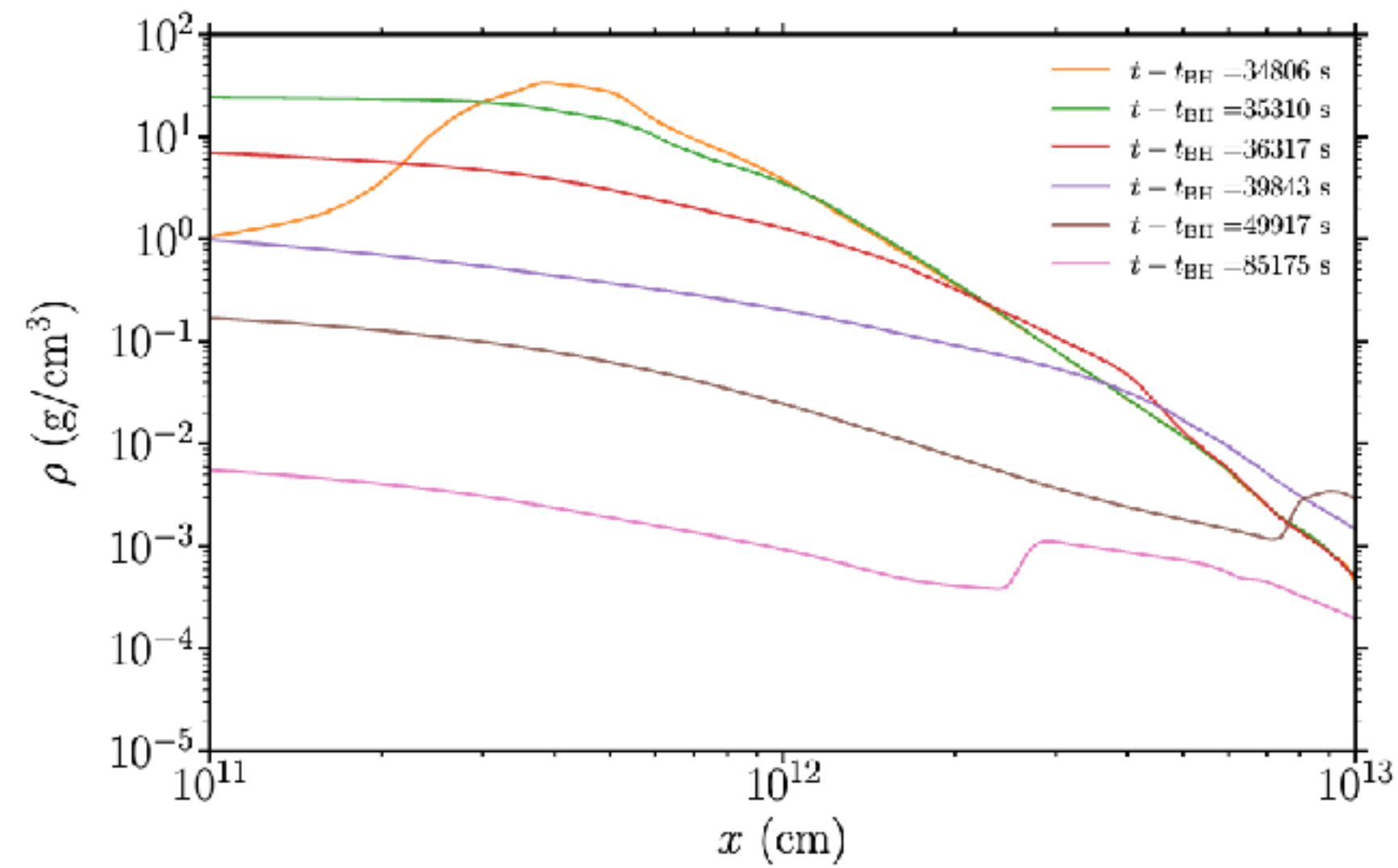


Result: Properties of ejecta



- Ejecta mass $\sim 1\%$ of initial SMS mass for fast-rotating SMS.
- Velocity $\sim 0.2 c$
- Kinetic E $\sim 10^{-4} M c^2$

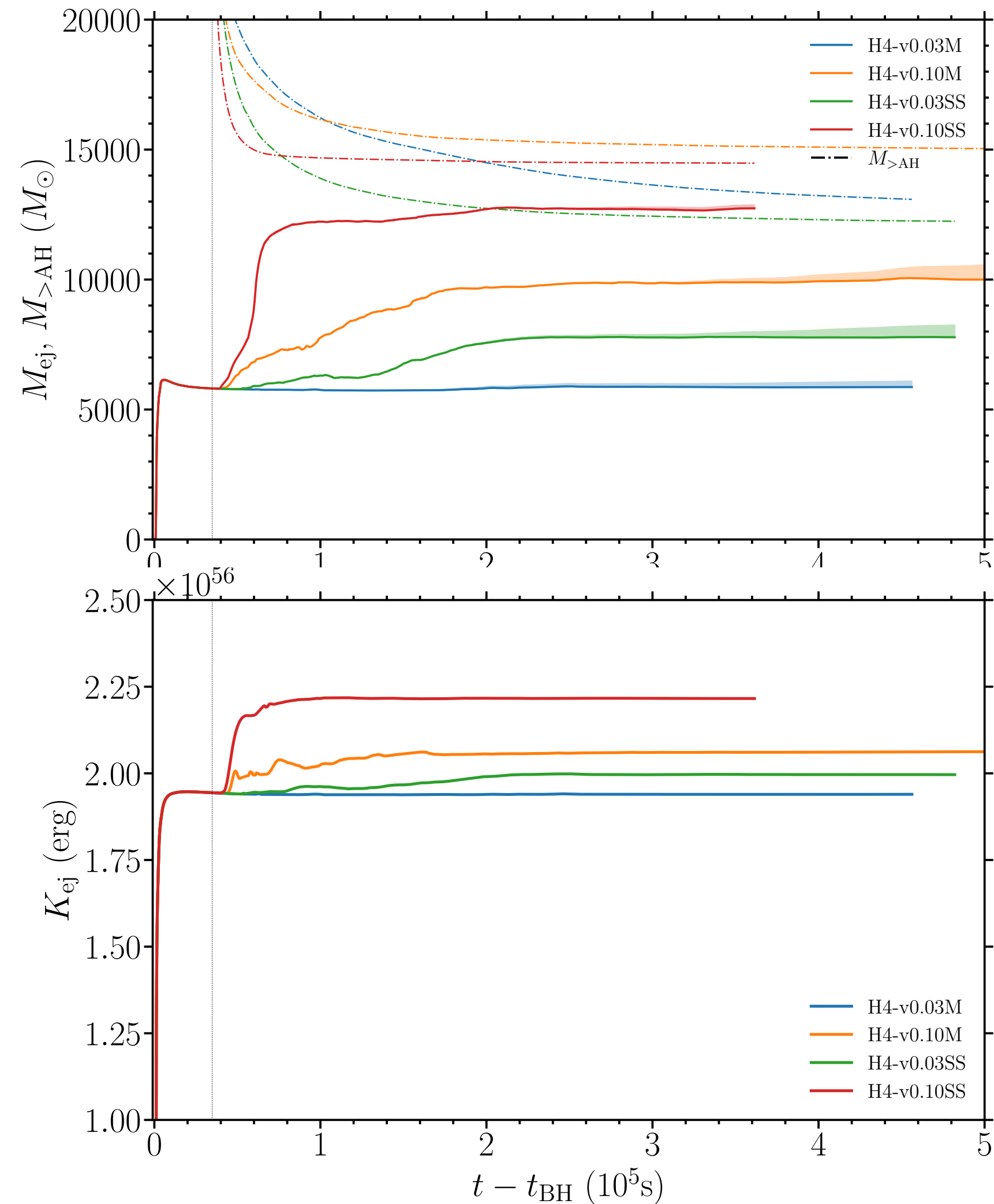
Result: Viscous evolution of the disk



Accretion timescale $\sim 10^4$ s

Up to $\sim 10M_{\odot}/\text{s} \sim 10^{13}\dot{M}_{\text{Edd}}$ (Hyper-Eddington)

Result: Viscous evolution of the disk

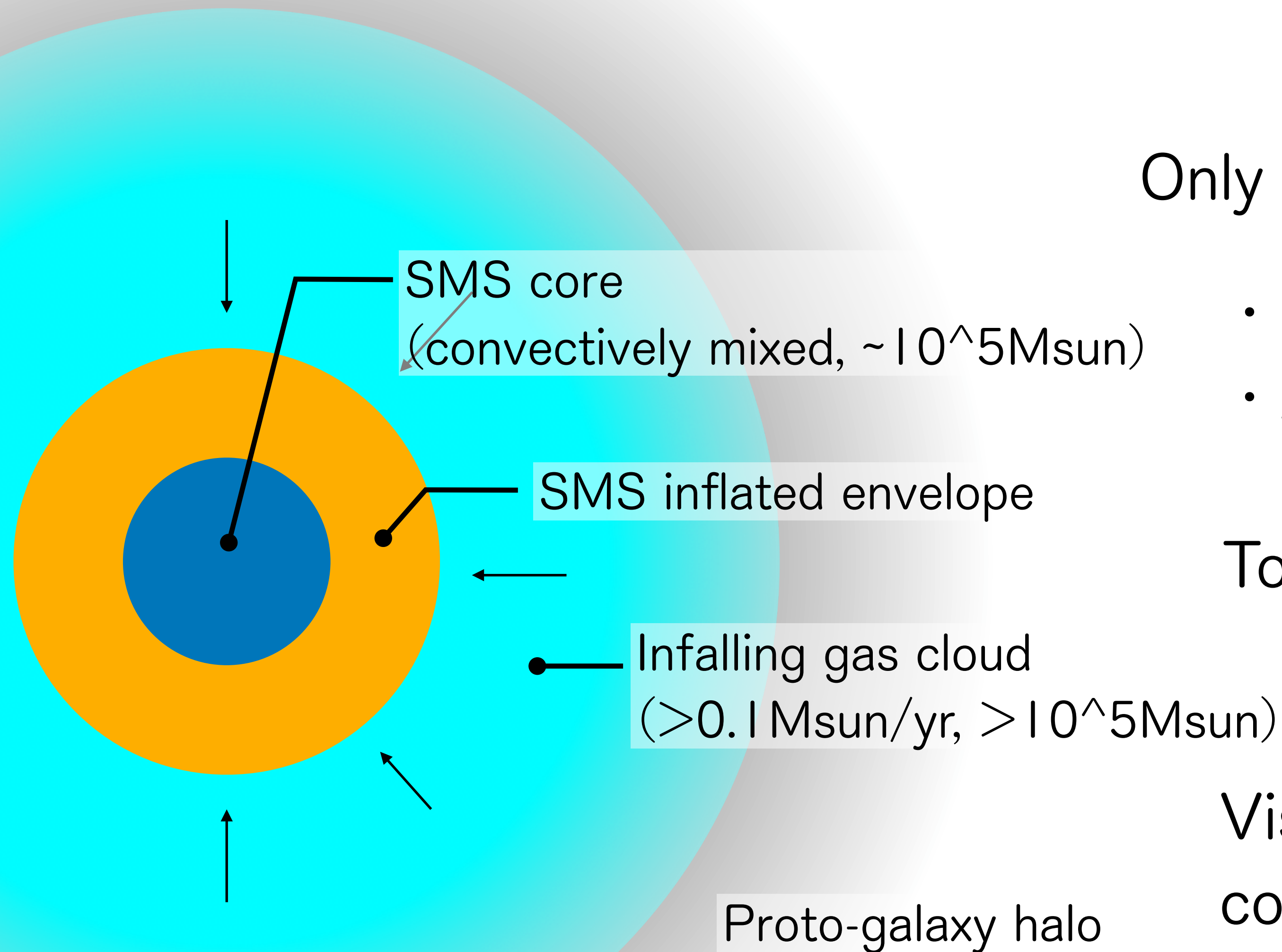


Viscosity-driven ejecta
(with different prescription & strength)

Ejecta mass can be $\sim 3 \times$ bounce-driven ejecta

Velocity $\sim 0.05 c \approx 1/4 \times$ bounce-driven ejecta
→ effect is minor in kinetic energy

Discussion: Realistic environment



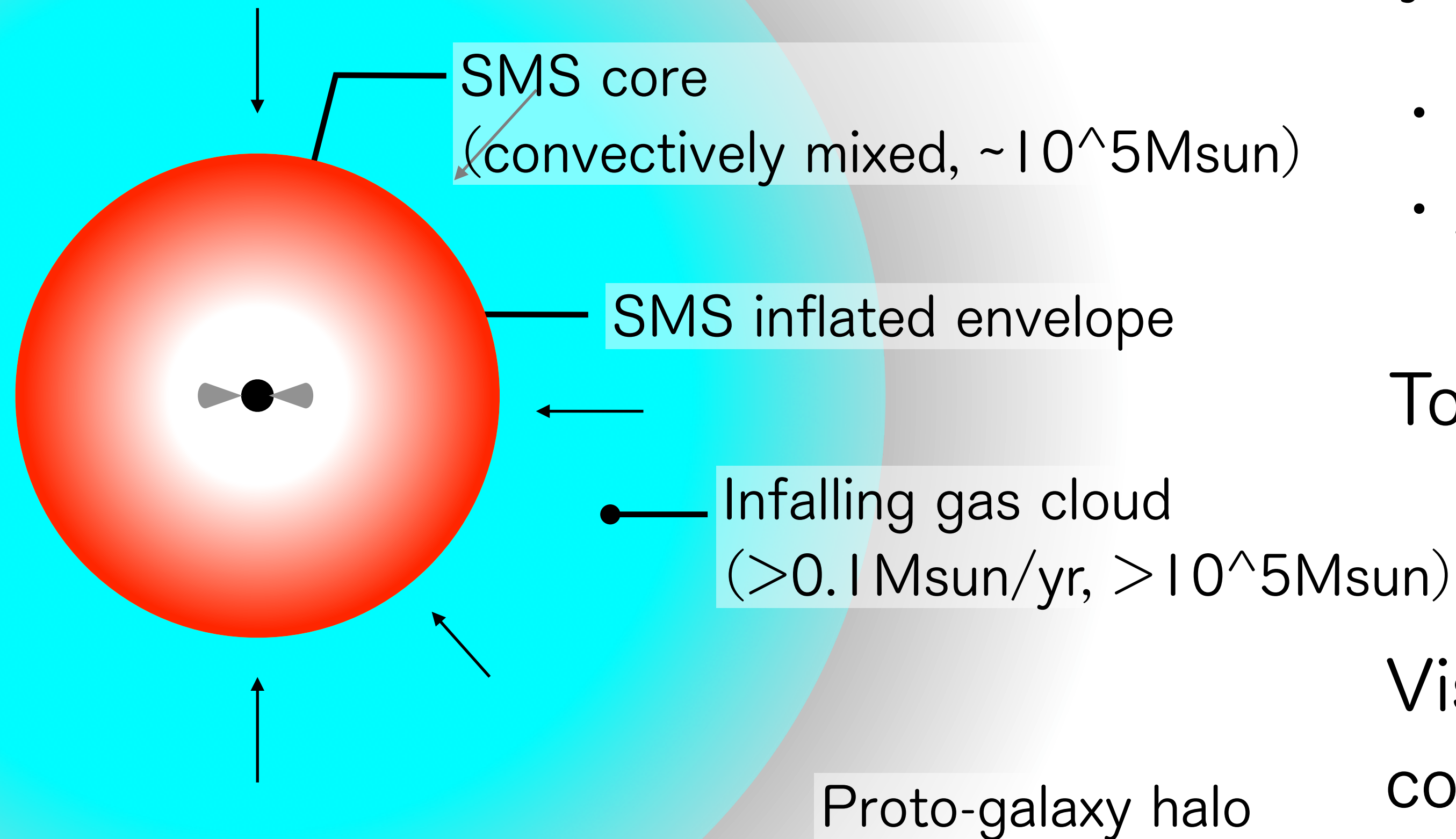
Only the collapse of SMS core is simulated.

- Envelope
- Atomic cooling cloud \sim SMS mass

Total ejecta mass $\sim 10^5 M_{\text{sun}}$

Viscosity-driven ejecta does not contribute much to total ejecta property

Discussion: Realistic environment



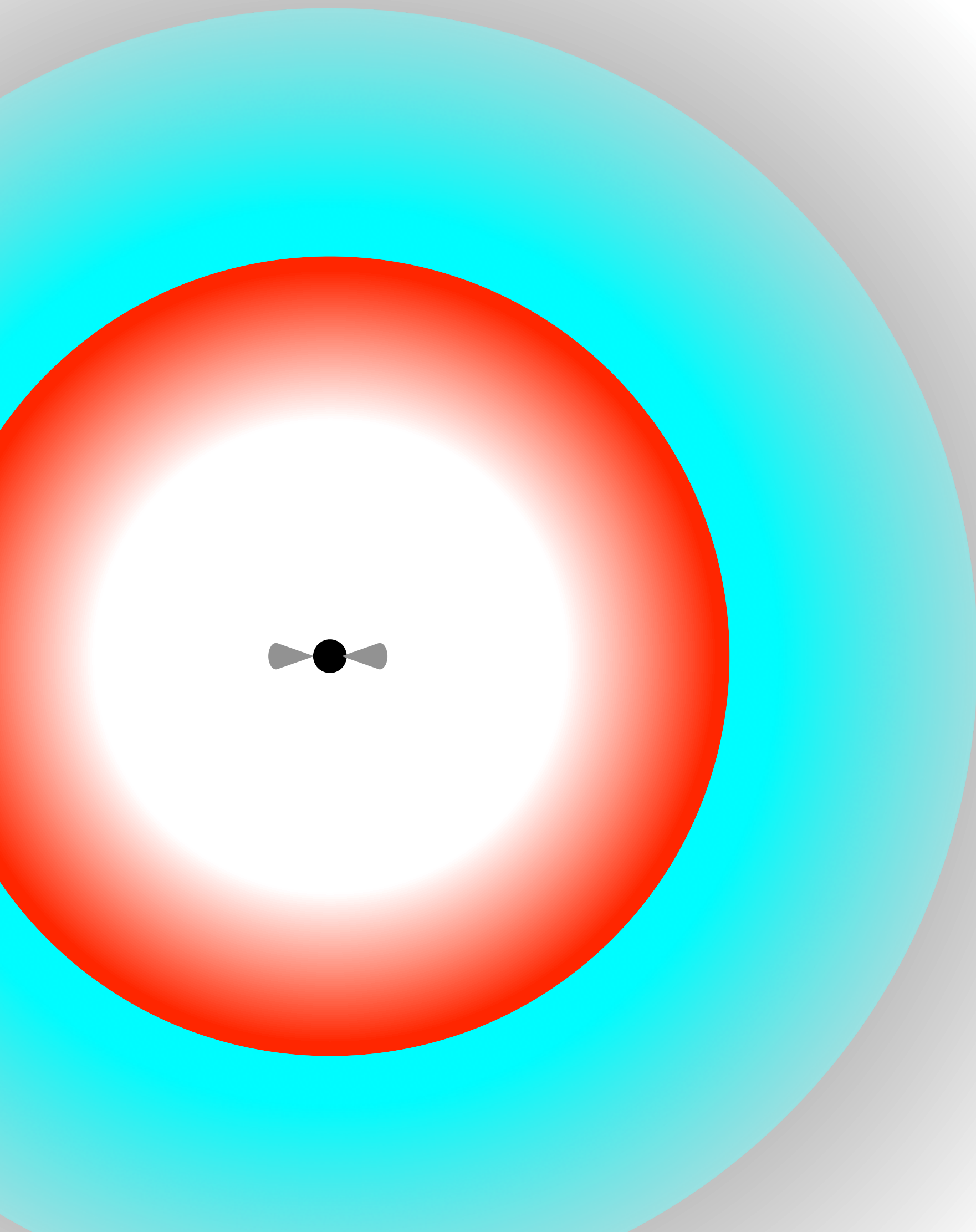
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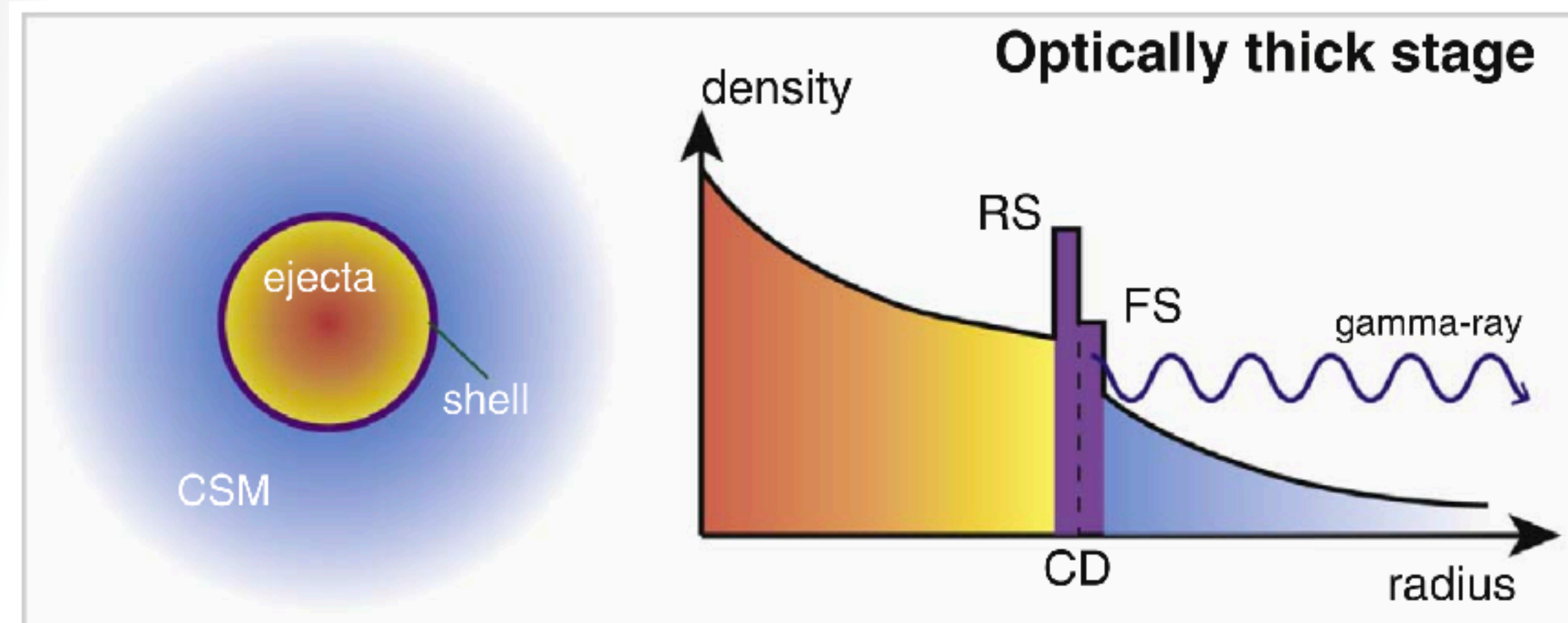
The ejecta sweep up inflated envelope

→ break out from SMS surface

The ejecta sweep up the infalling gas cloud

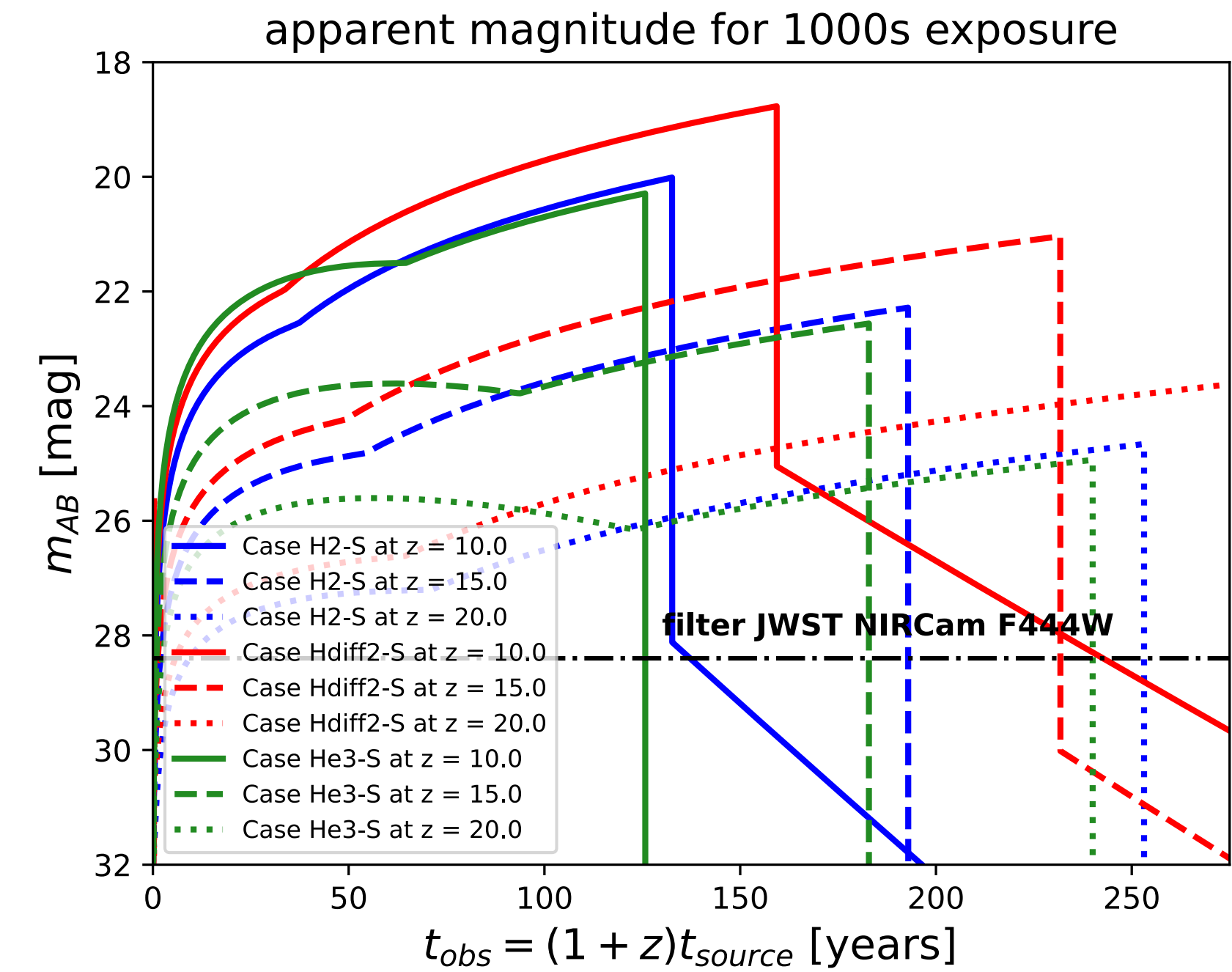
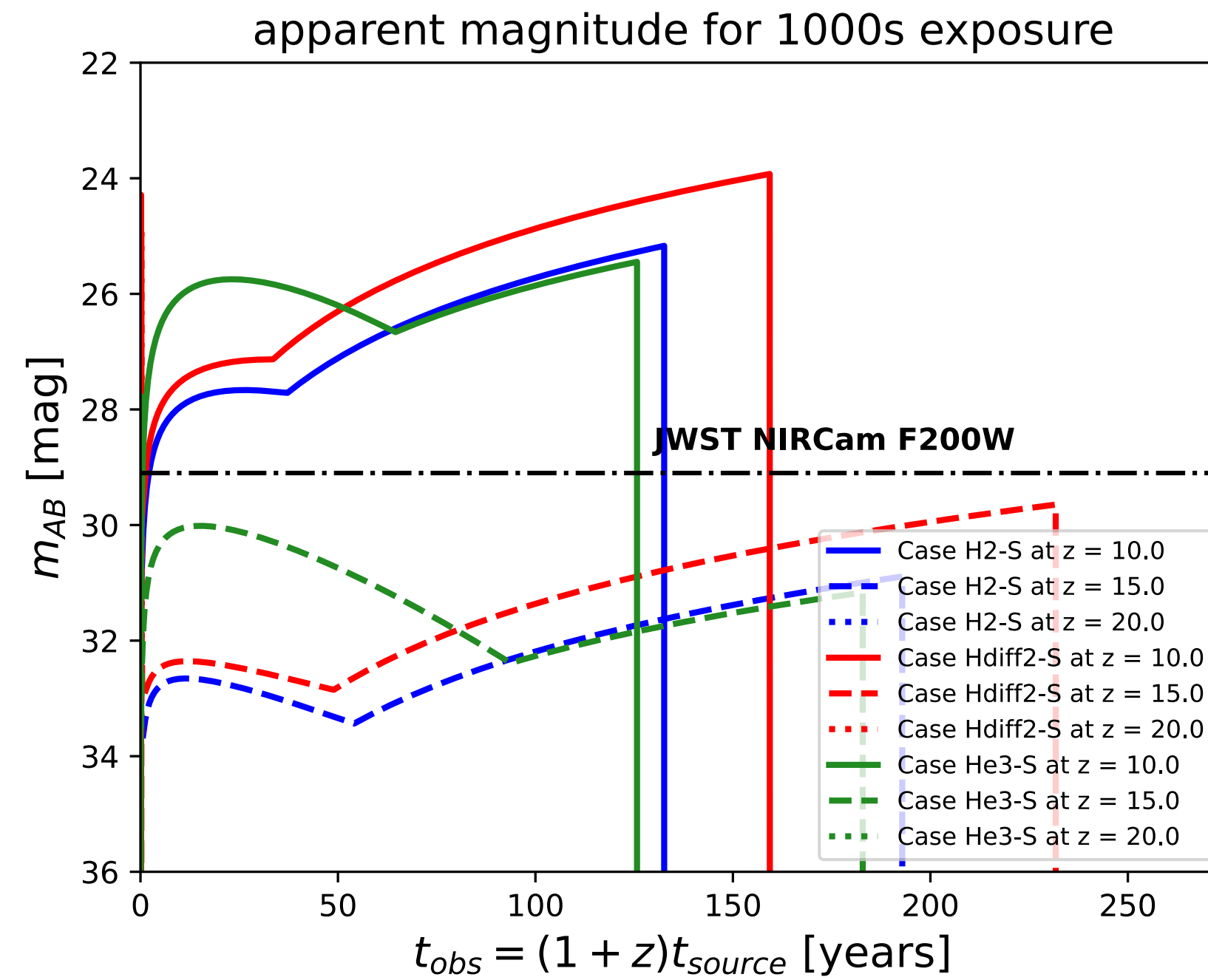
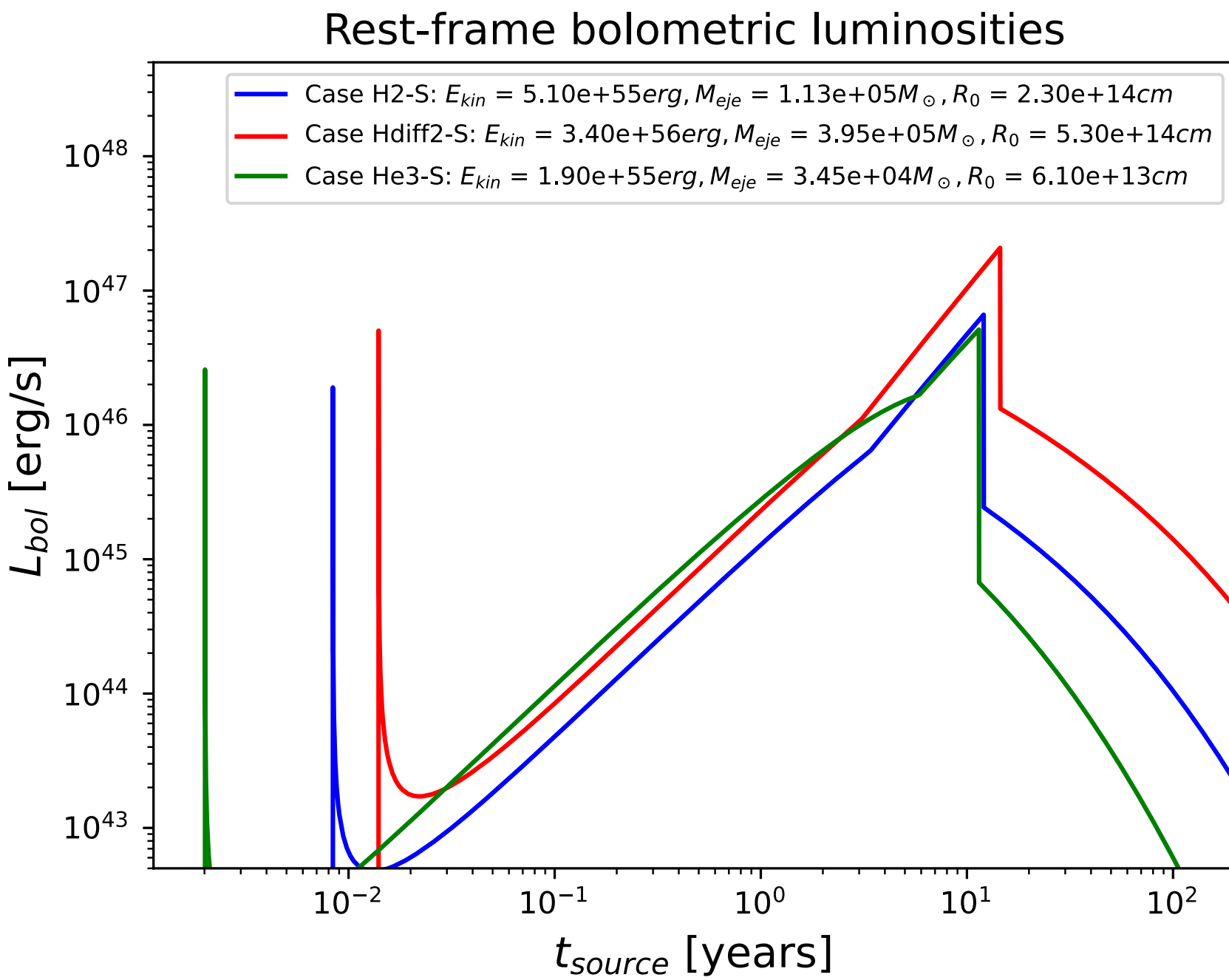
~ CSM-interacting SN

(with $E \sim 10^{55} - 10^{56}$ erg, $M_{\text{CSM}} \sim 10^5 M_{\text{sun}}$)

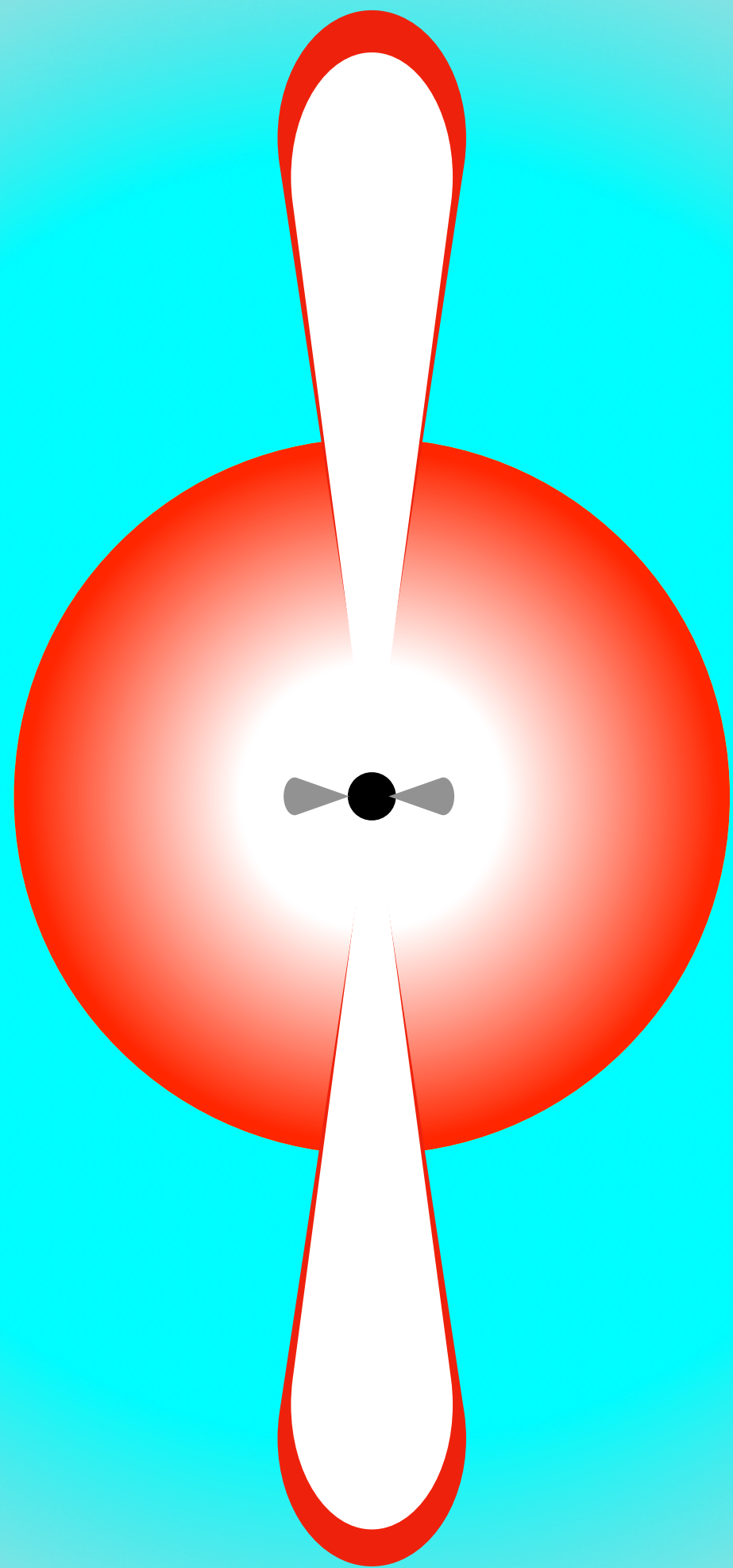


Observational feature

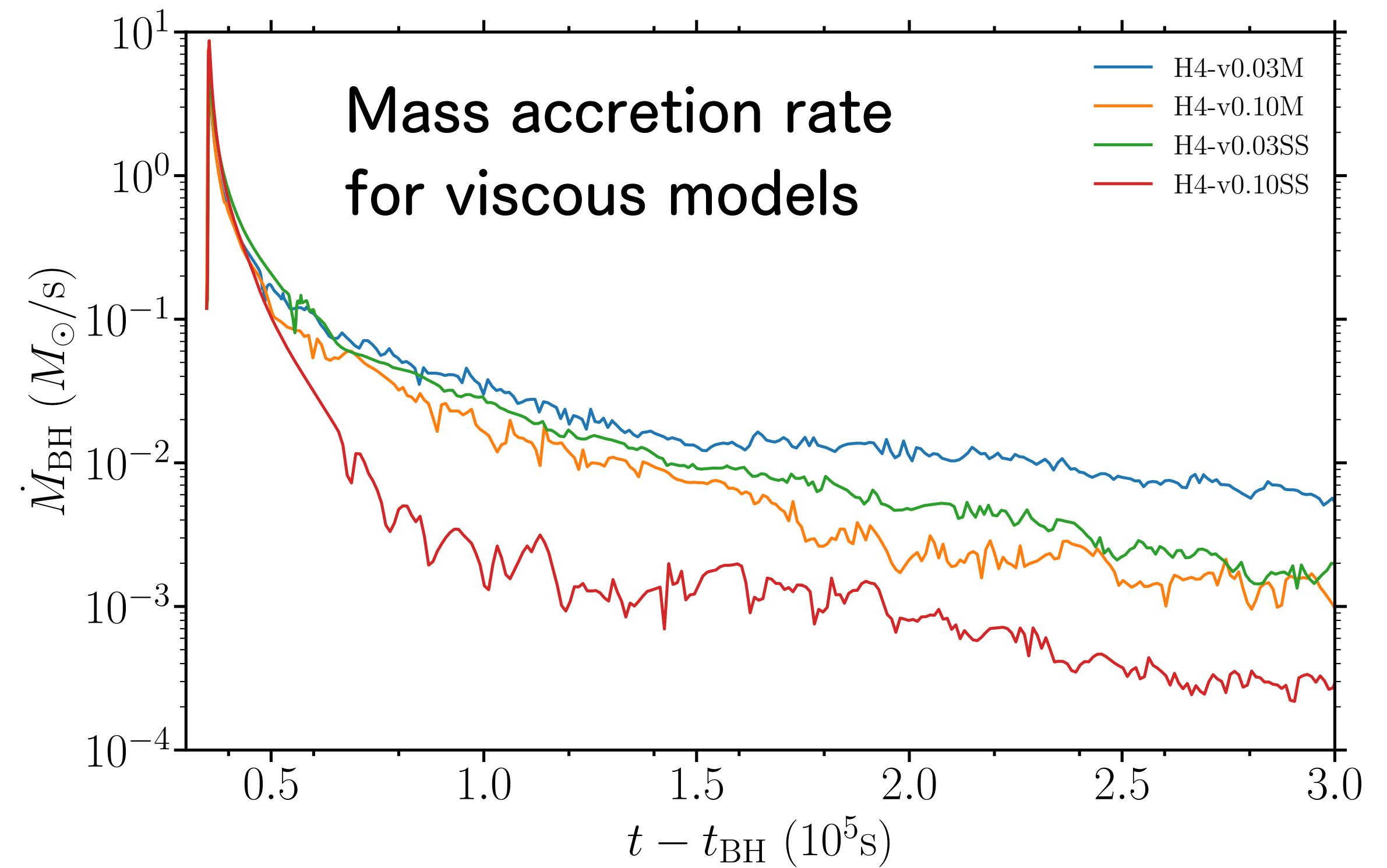
Jockel, SF+ in prep.



Discussion: Jet driven by BH-disk



Disk accretion onto BH could drive relativistic jet cf. Matsumoto+15,16



Summary (Part I)

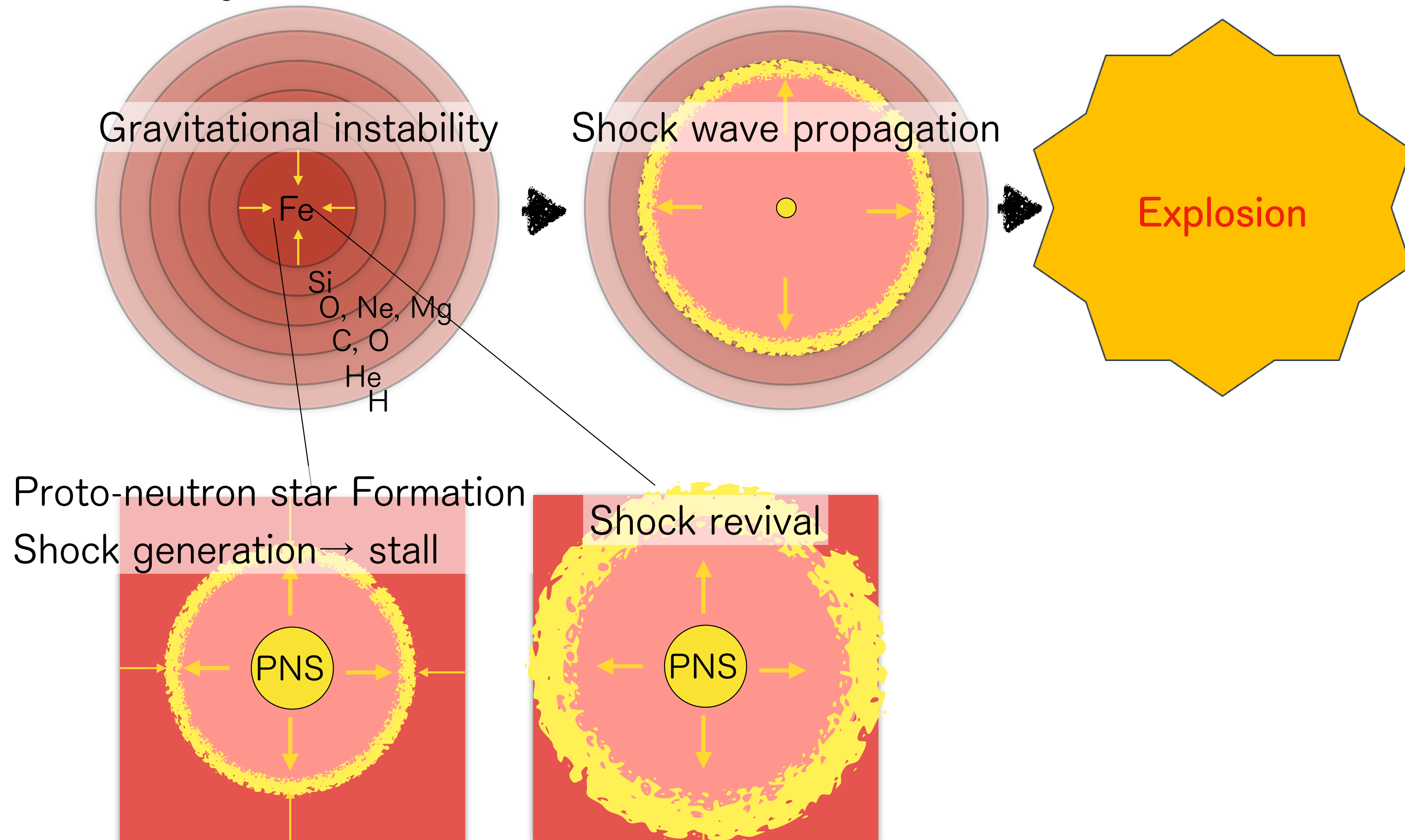
- ✓ Numerical-relativity simulations of the collapses of rotating supermassive stars
- ✓ Bounce-shock-induced ejecta up to 1% of core mass ($10^{-2}M_{\text{core}}$), $v \approx 0.2c$
(Mass is likely dominated by the swept-up cloud surrounding the star)
- ✓ Kinetic energy $\sim 10^{55} - 10^{56}$ erg ($10^{-4}M_{\text{core}}c^2$)
- ✓ Bright CSM-interacting SNe with duration $\sim 10(1+z)$ yr

Outline

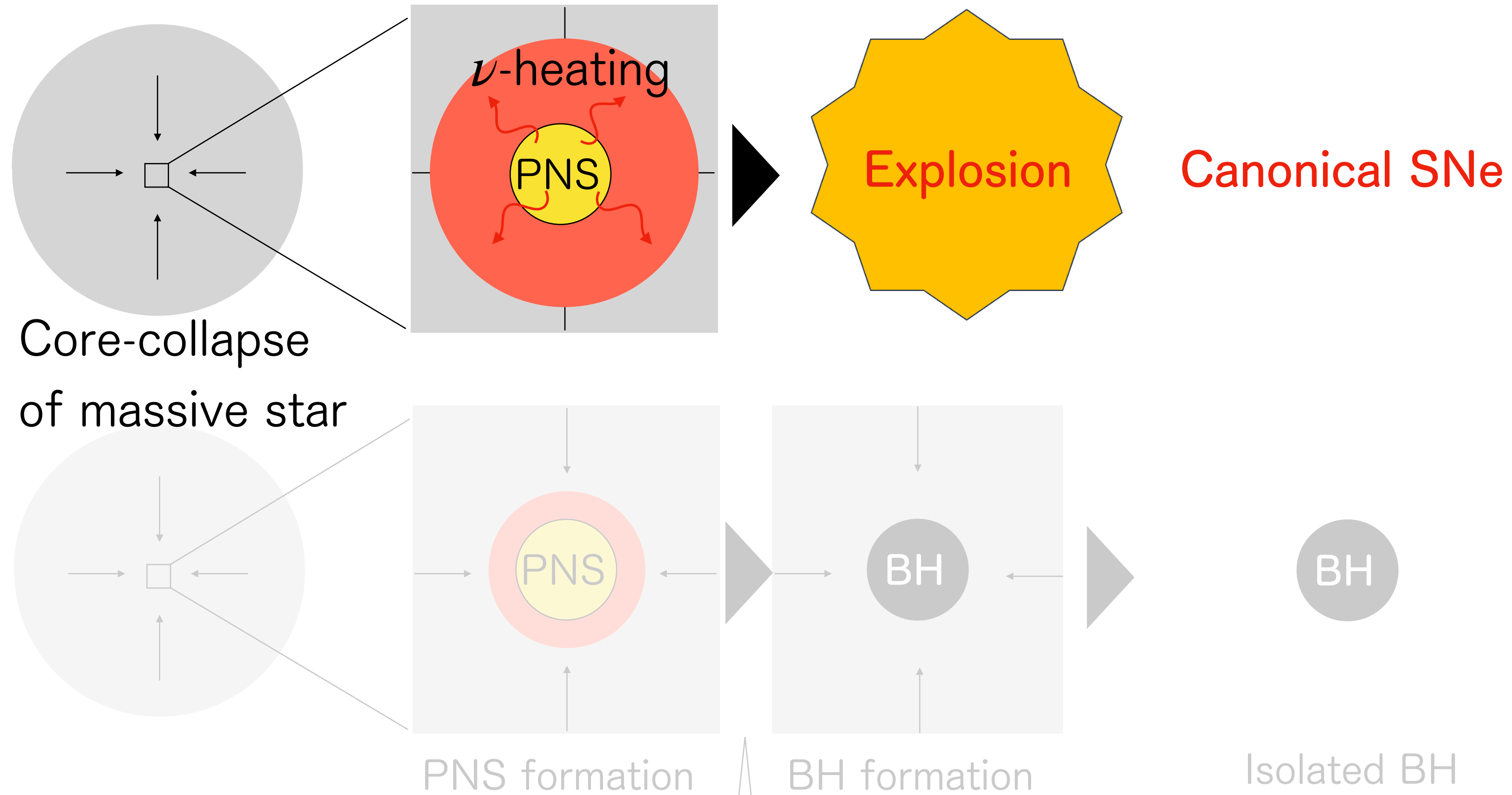
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Core-collapse supernovae

Stars with mass $\gtrsim 10 M_{\odot}$

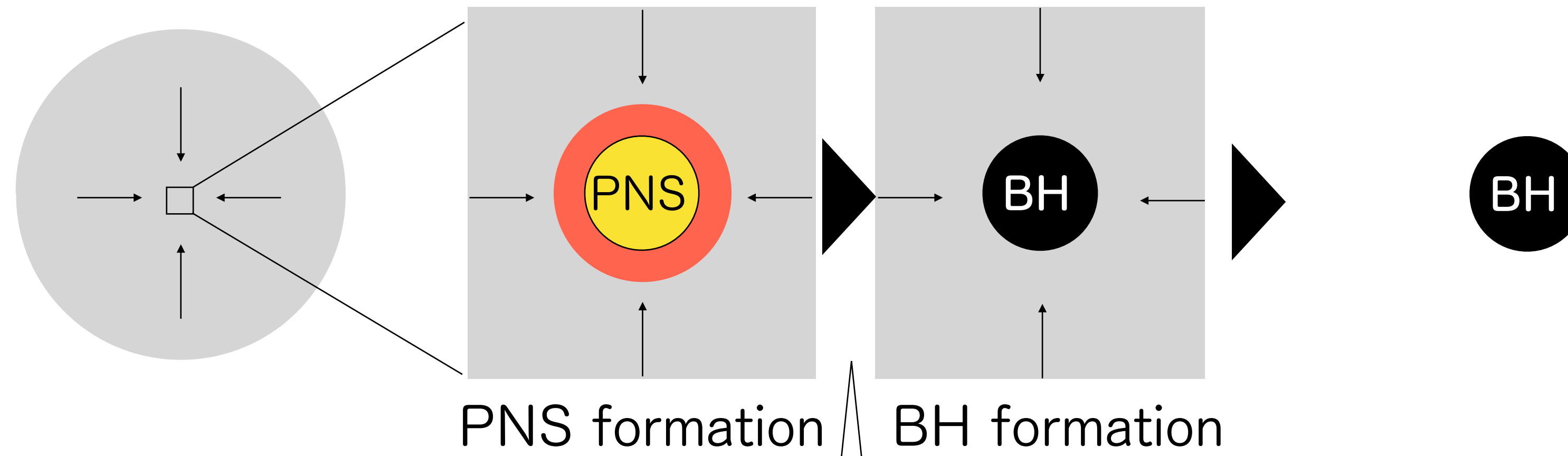
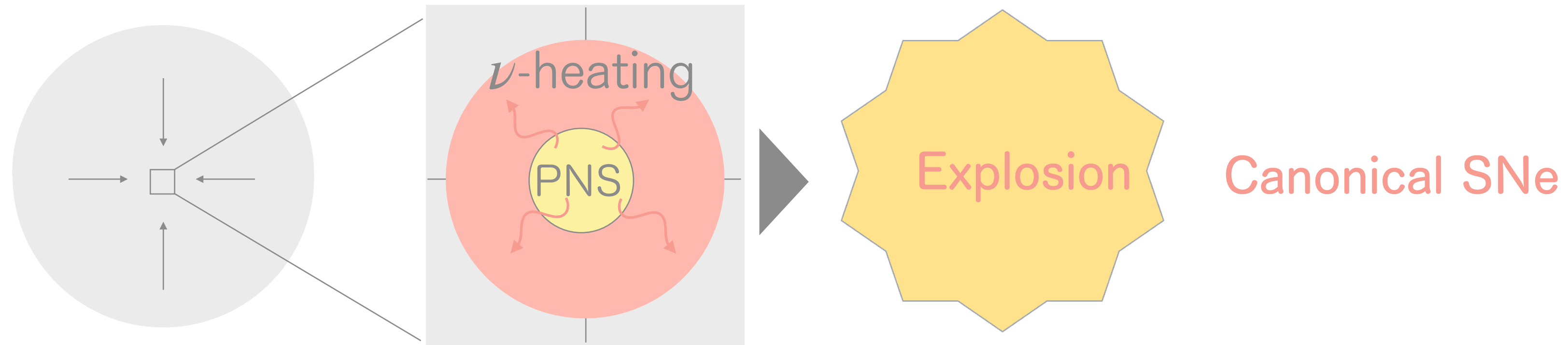


Collapsar



Explosion during PNS phase fails if, e.g., the core compactness is too high.

Collapsar

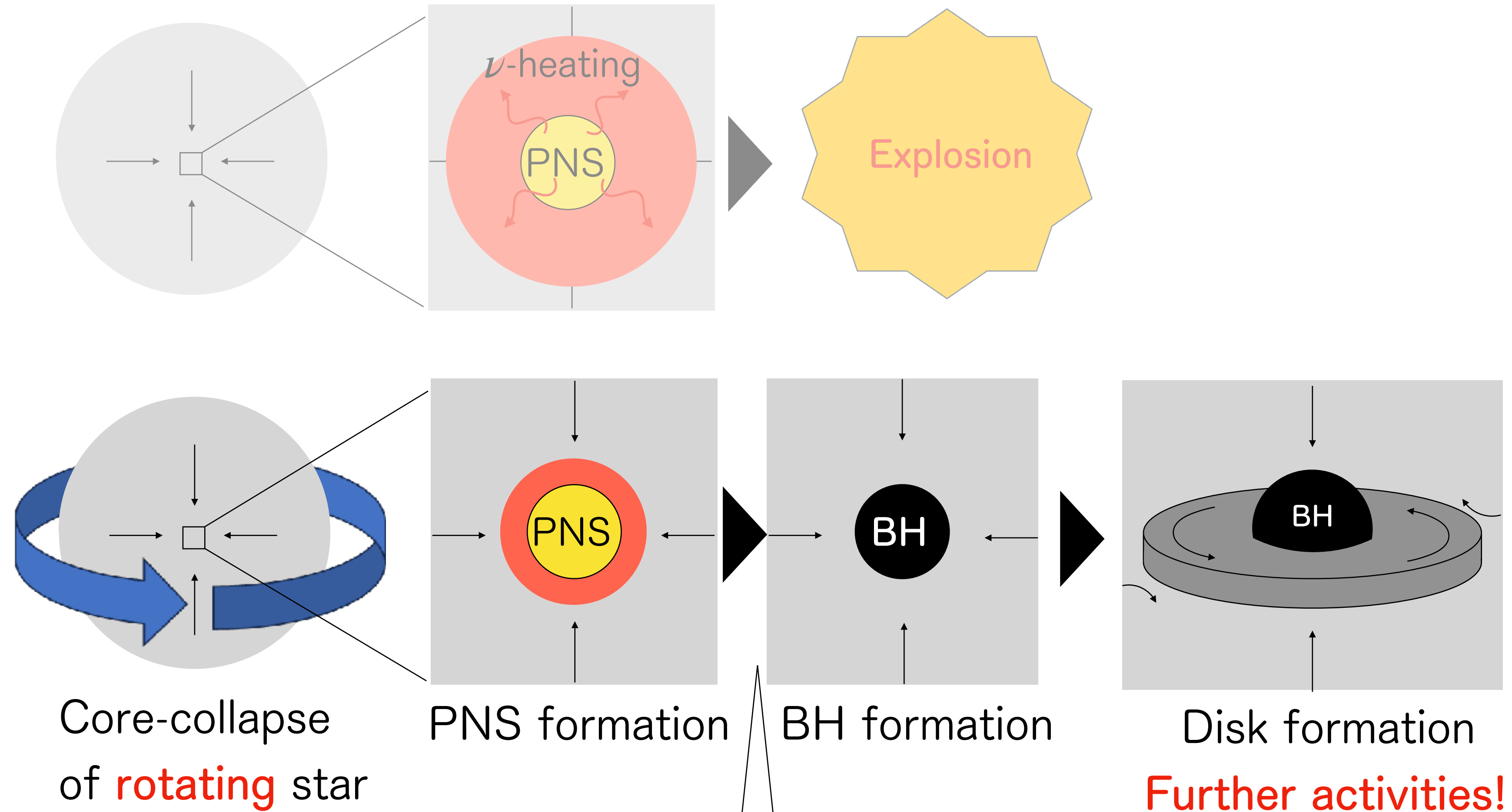


Explosion during PNS phase fails if, e.g., the core compactness is too high.

Note: MHD process can help explosion

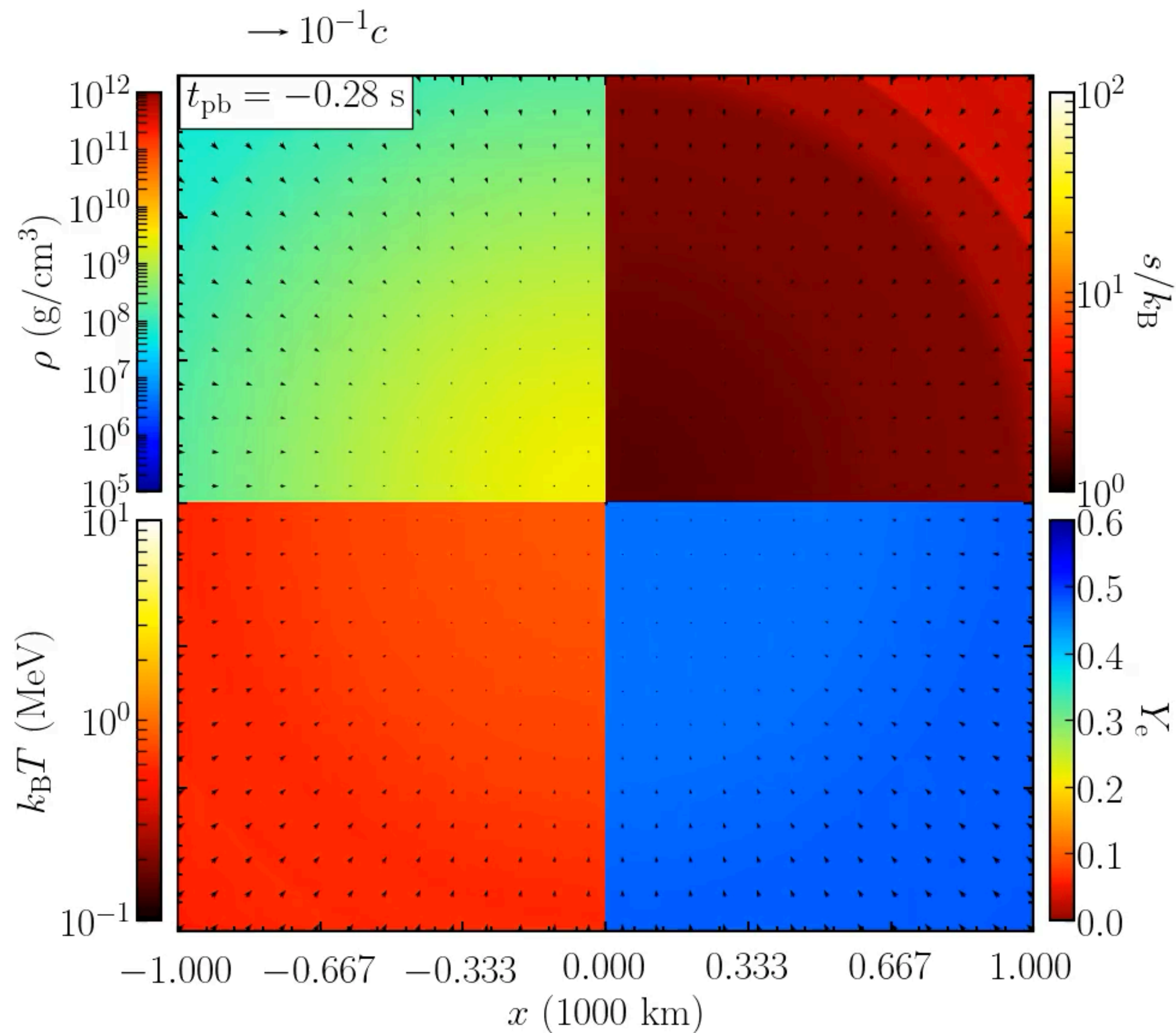
e.g., Obergaulinger & Aloy

Collapsar



Explosion during PNS phase fails if, e.g., the core compactness is too high.

Collapsar



An example SF+2023

Collapse of $M_{\text{ZAMS}} = 20M_{\odot}$ star

(taken from Aguilera-Dena+2020)

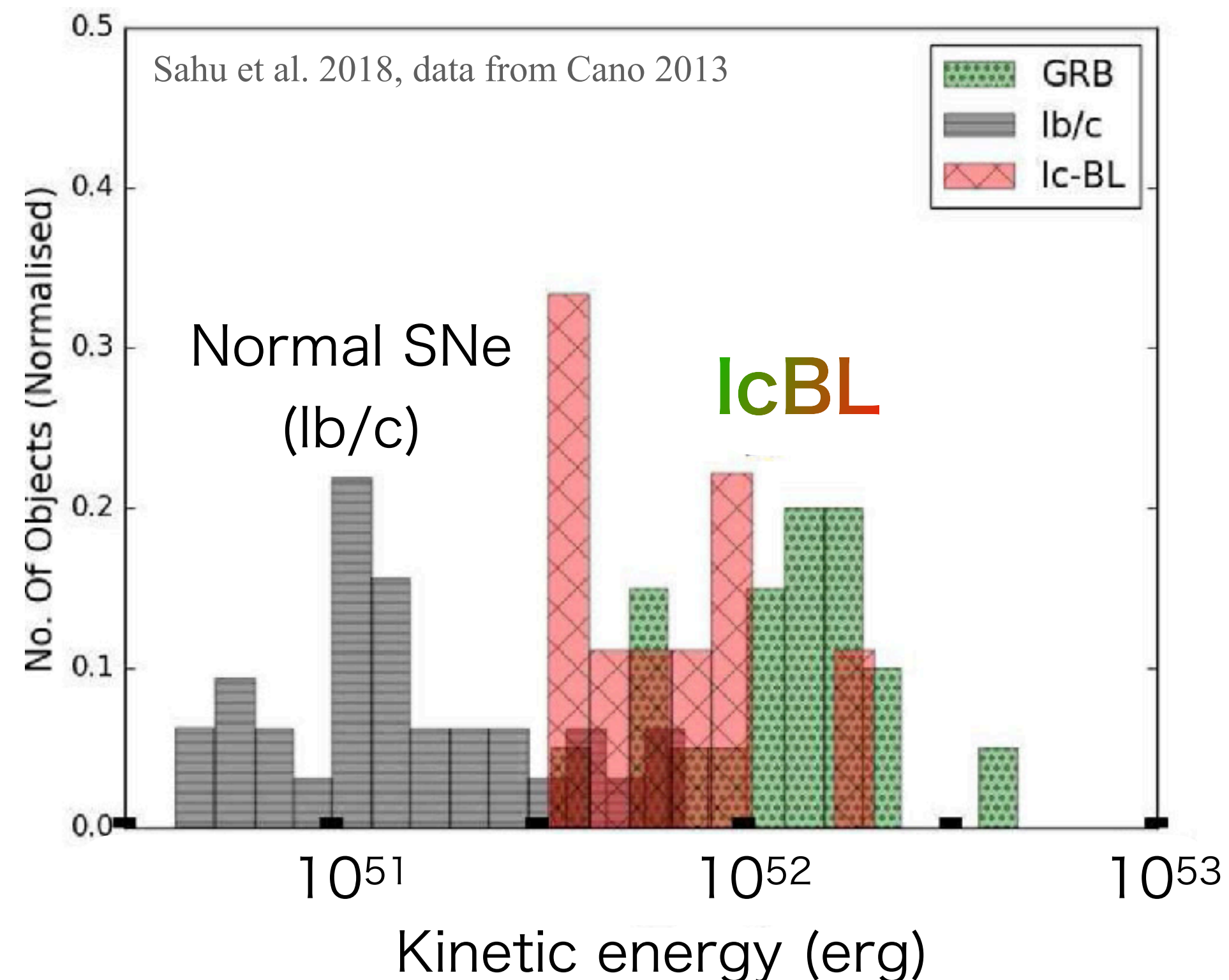
BH-disk activities and GRB-SN

Gamma-ray bursts (GRBs)

BH-disk is one of the promising central engines (e.g., Woosley et al. 1993...)

Broad-lined type Ic SNe (SNe Ic-BL; Hypernovae)

Long GRBs are accompanied by energetic SNe (Ic-BL)



- Explosion (kinetic) energy $E_K = (0.8 - 4.4) \times 10^{52}$ erg
- ^{56}Ni mass $M_{\text{Ni}} = (0.2 - 0.5) M_{\odot}$ (Cano et al. 17)

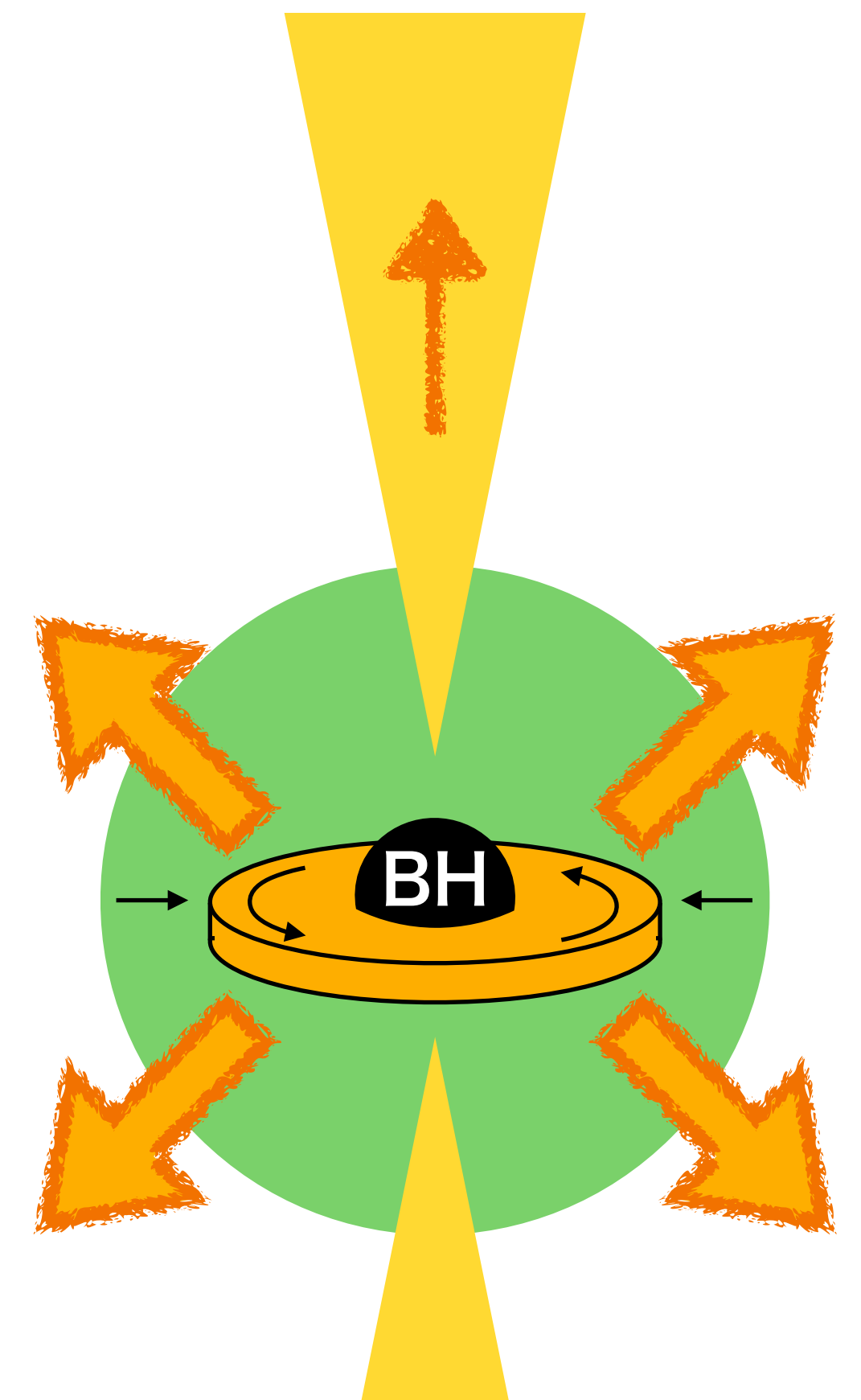
BH-disk activities and GRB-SN

Disk outflow (MacFadyen & Woosley 1999)

Energy generated by viscous accretion:

$$\frac{GM_{\text{BH}}M_{\text{disk}}}{r_{\text{disk}}} \approx 3 \times 10^{52} \text{ erg} \left(\frac{M_{\text{BH}}}{10M_{\odot}} \right) \left(\frac{M_{\text{disk}}}{0.1M_{\odot}} \right) \left(\frac{r_{\text{disk}}}{10^7 \text{ cm}} \right)^{-1}$$

Viscosity-driven outflow from disk would naturally explain such SNe



Neutrino cooling vs viscous heating

- Neutrino emission cooling $t_{\text{weak}} \sim \frac{1}{G_{\text{F}}^2 T^5} \approx 1 \text{ s} \left(\frac{kT}{1 \text{ MeV}} \right)^{-5}$

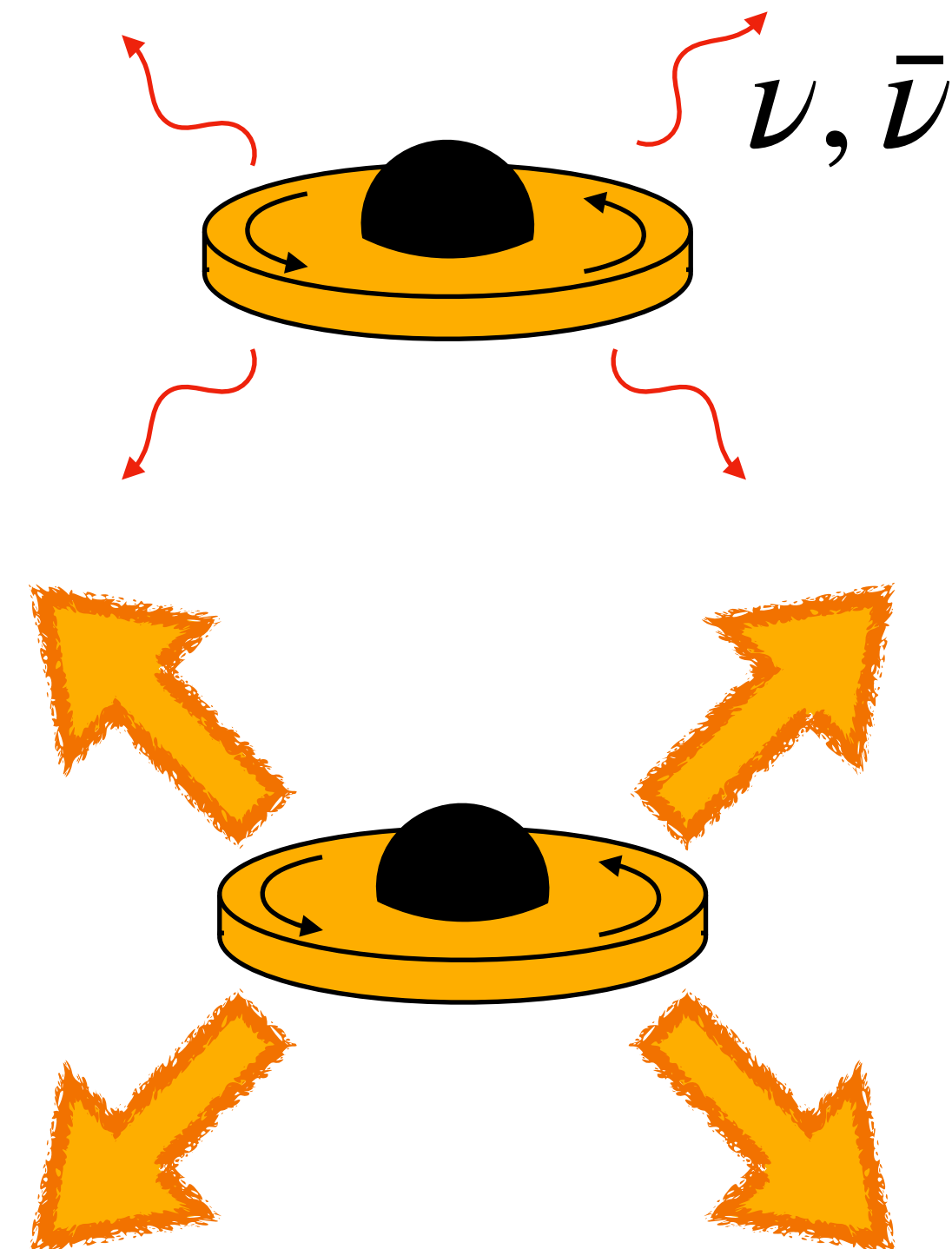
- MHD turbulence \rightarrow Viscous angular momentum transport/heating

$$t_{\text{vis}} \sim \frac{R^2}{\nu} = 1 \text{ s} \left(\frac{R}{10^7 \text{ cm}} \right) \left(\frac{c_s}{10^9 \text{ cm/s}} \right)^{-1} \left(\frac{\alpha}{0.03} \right)^{-1} \left(\frac{H/R}{0.3} \right)^{-1}$$

✓ $t_{\text{weak}} \lesssim t_{\text{vis}}$ (NDAF) phase: weak/no outflow

✓ $t_{\text{weak}} \gg t_{\text{vis}}$ phase: viscosity can drive outflow

Same as NS-merger remnant disk

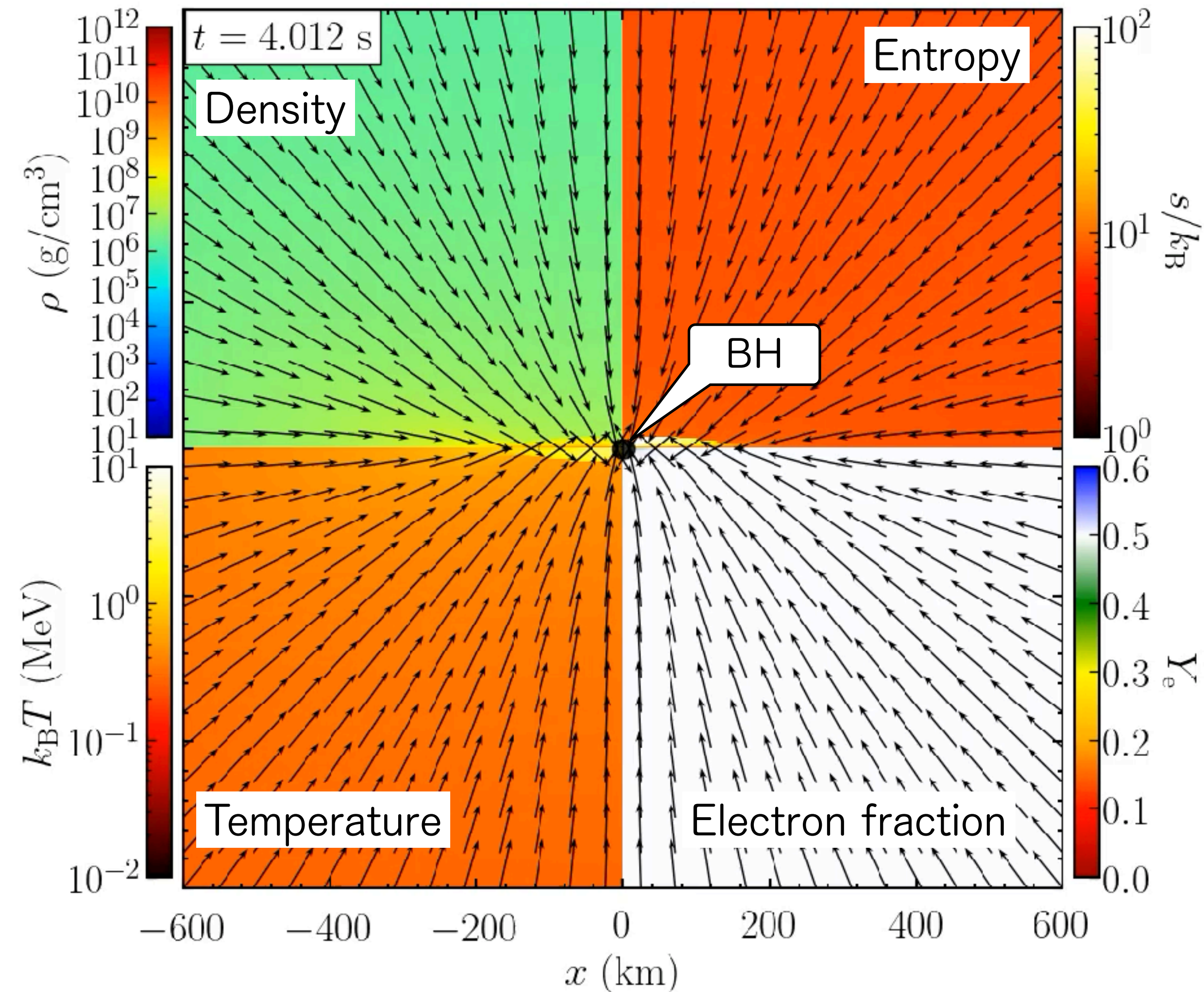


2D-axisymmetric simulation with solving

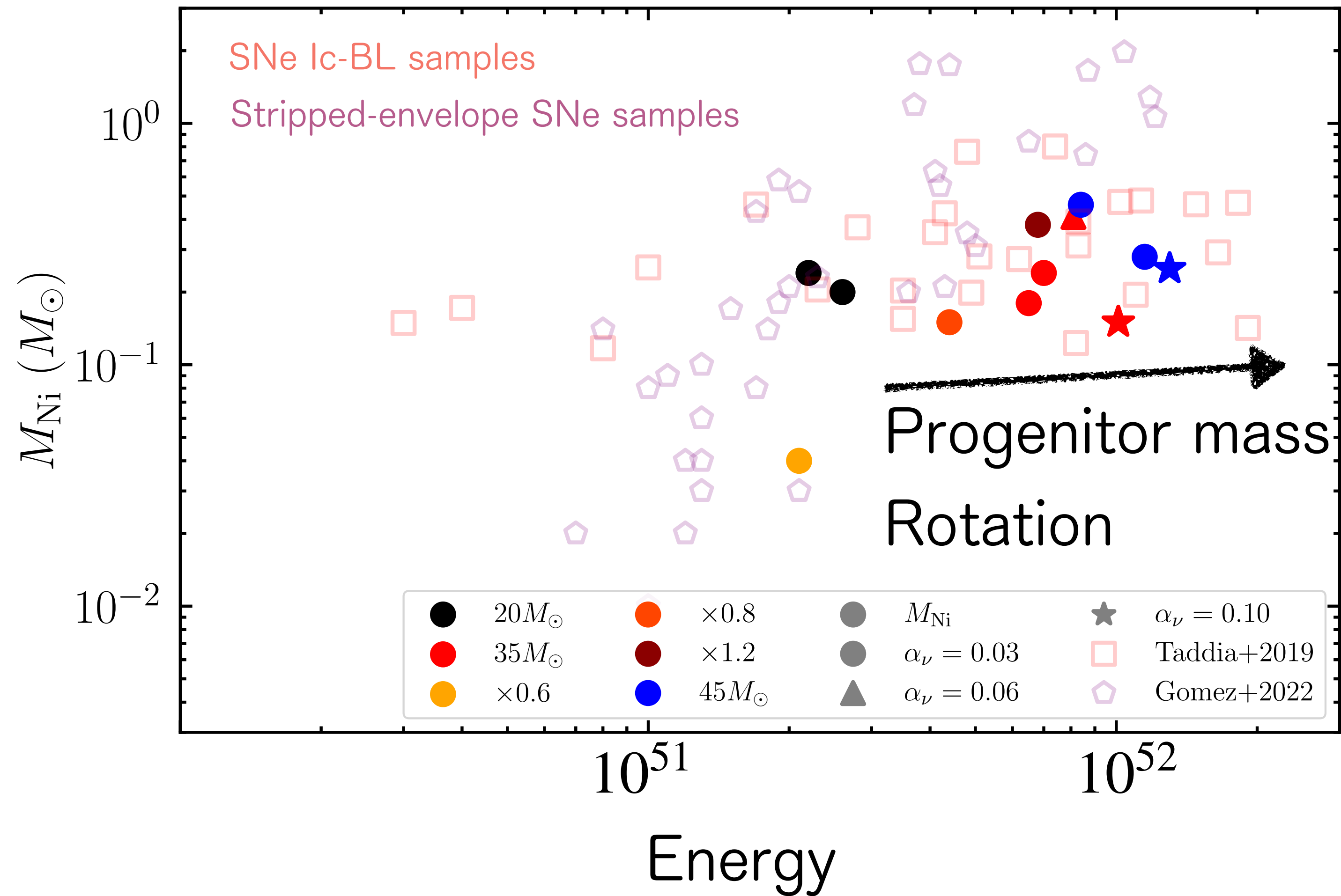
- ✓ Einstein's equation Nakamura & Shibata 95, Baumgarte & Shapiro 99
- ✓ Neutrino radiation transfer equation Thorne 81, Shibata et al. 11
- ✓ Viscous hydrodynamics equation Israel & Stuart 79, Shibata et al. 17, Shibata & Kiuchi 17
(to mimic MHD turbulence)

Disk formation \rightarrow NDAF \rightarrow Outflow

Progenitor: $M_{\text{ZAMS}} = 35M_{\odot}$ star



Comparison with observations



Nucleosynthesis calculation in the ejecta $\rightarrow M_{\text{Ni}} \gtrsim 0.1M_{\odot}$

BH-disk can power the energetic explosion.

MHD models for GRB jets

Only with viscosity, jet cannot be produced.

With MHD, we have Blandford-Znajek (BZ) process Blandford & Znajek (1977)

$$L_{\text{BZ}} \sim (BM\chi)^2 \sim 10^{50} \text{erg/s} \left(\frac{\chi}{0.7}\right)^2 \left(\frac{M}{10M_{\odot}}\right)^2 \left(\frac{B}{10^{14}\text{G}}\right)^2$$

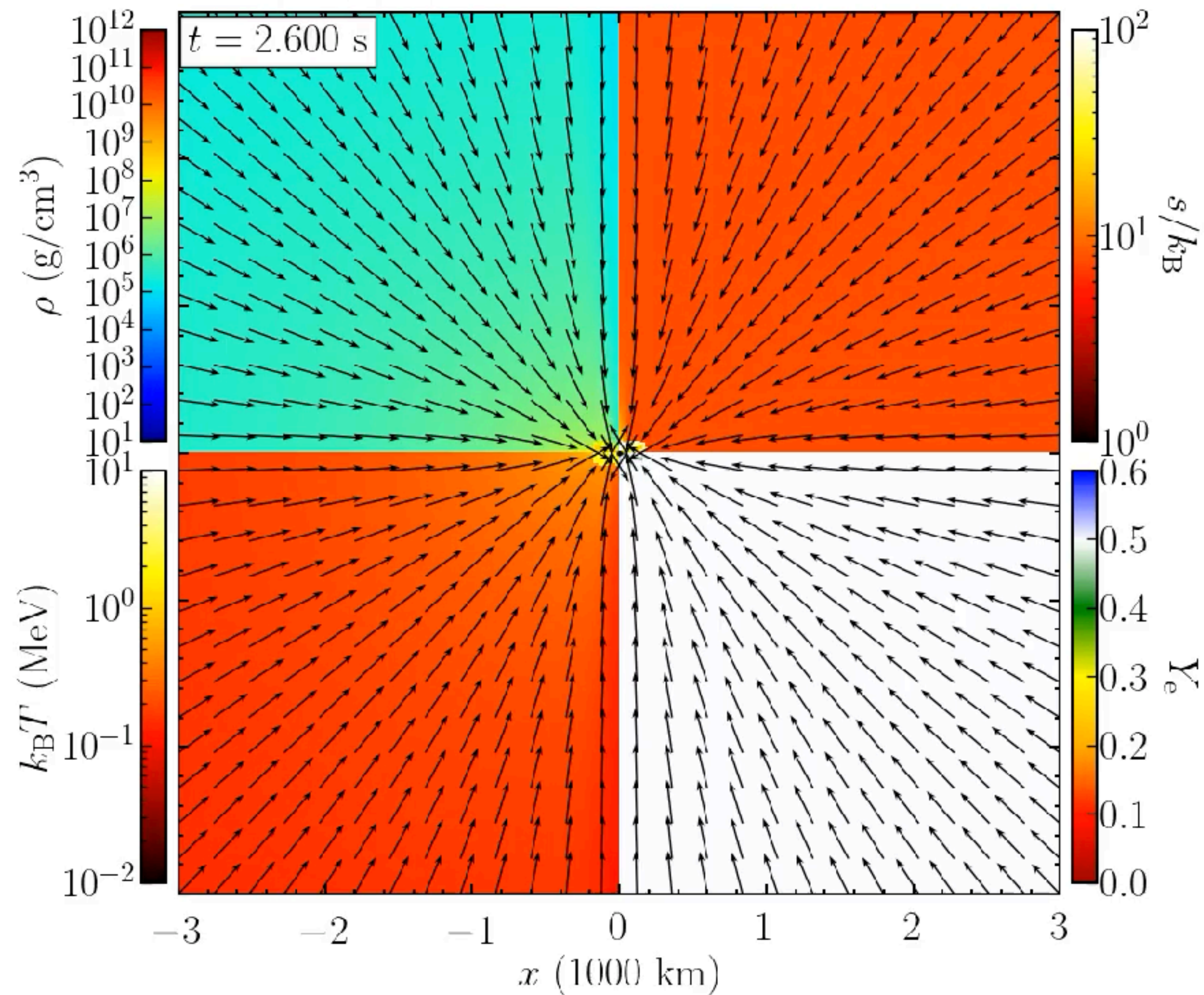
2D-axisymmetric simulation with solving

- ✓ Einstein's equation Nakamura & Shibata 95, Baumgarte & Shapiro 99
- ✓ Neutrino radiation transfer equation Thorne 81, Shibata et al. 11
- ✓ Magneto-hydrodynamics equation Shibata, SF+21

MHD models for GRB jets

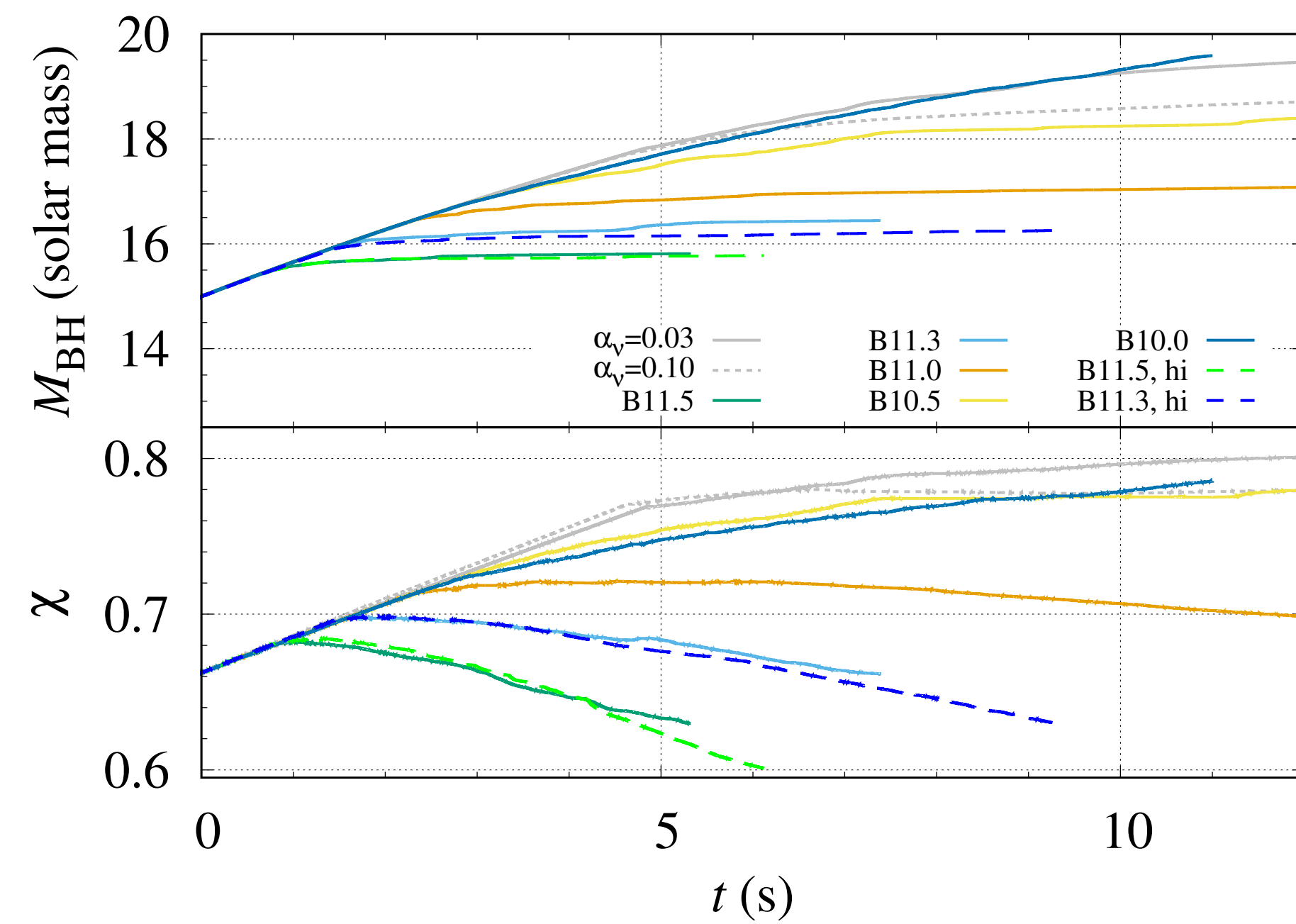
Shibata, SF+24

Progenitor: $M_{\text{ZAMS}} = 35M_{\odot}$ star, Poloidal field



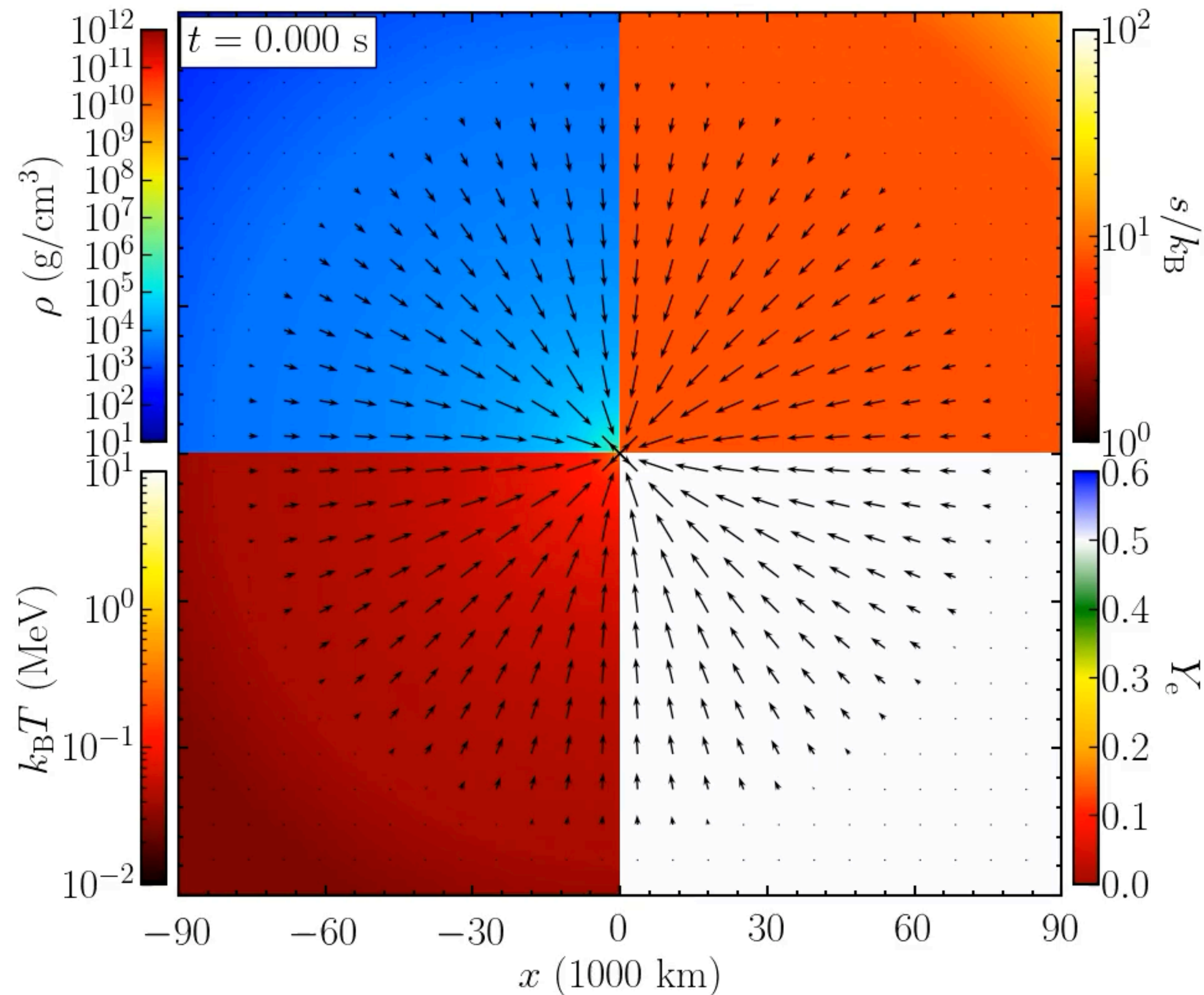
Fossil field \rightarrow BZ jets

Feedback on BH spin is numerically observed



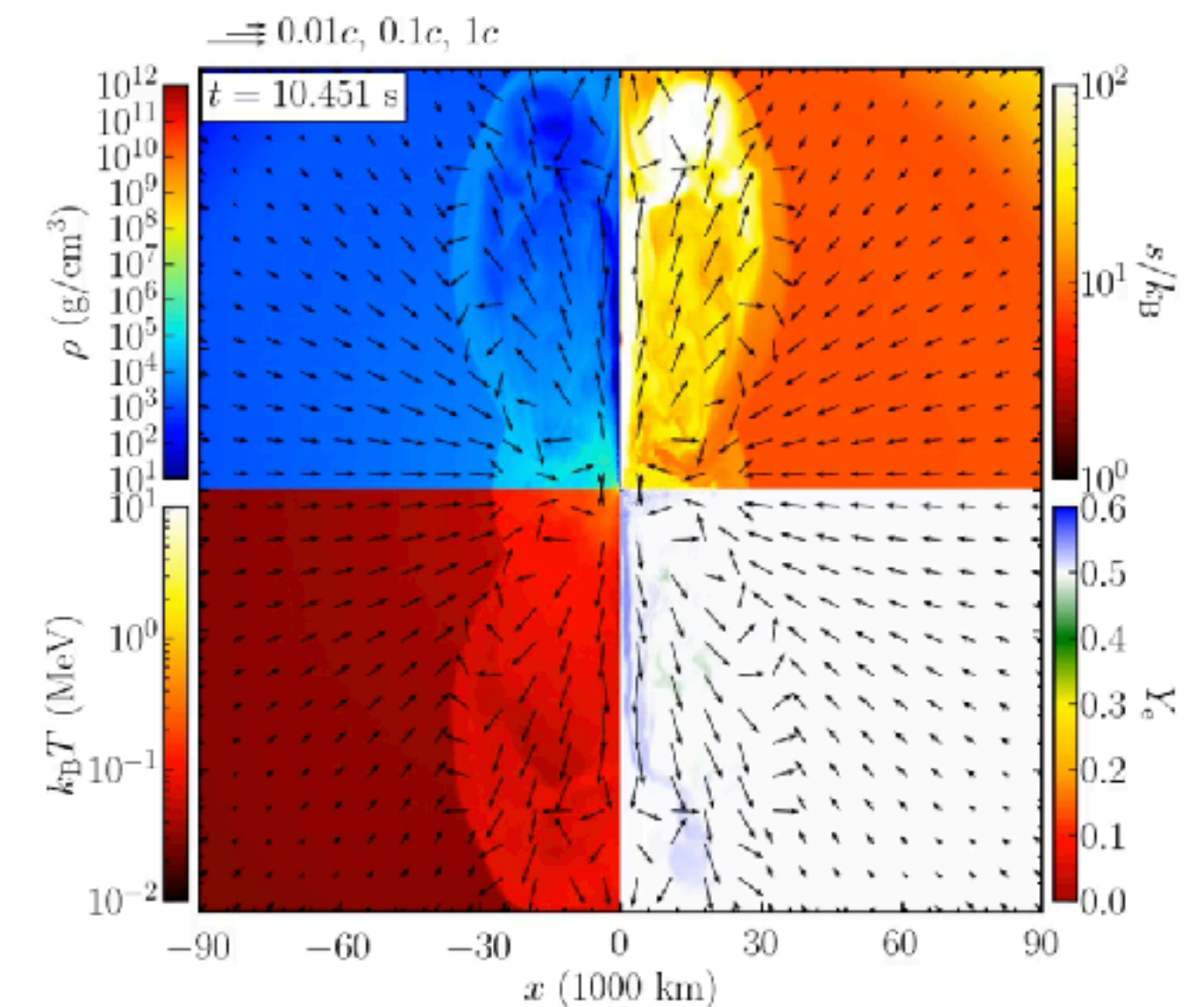
Collapsar with MHD+dynamo

Progenitor: $M_{\text{ZAMS}} = 35M_{\odot}$ star, toroidal field



Jet + spherical components

cf. Ideal MHD case



Summary (Part II)

Numerical relativity simulations of collapses of rotating massive stars

- ✓ It can explode with $E \sim 10^{52}$ erg driven by disk outflow (←viscous model)
- ✓ It can synthesize sufficient amount ($\gtrsim 0.1M_{\odot}$) of ^{56}Ni (←viscous model)
- ✓ It can drive a jet (←MHD model with an ideal config.)
- ✓ MHD+phenomenological dynamo model will come soon.
- ✓ No significant r-process in the ejecta